

Kinematics of Gravity–Capillary Waves Above an Evolving Underwater Current

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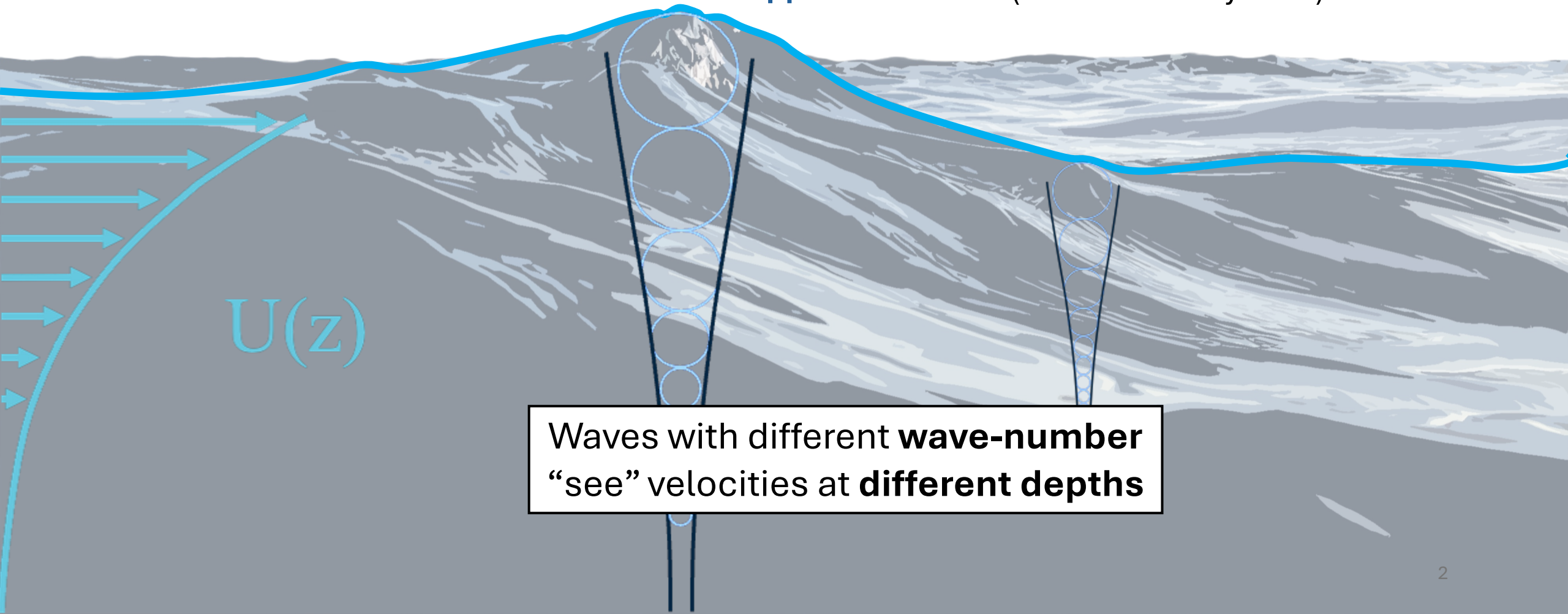
Introduction

$$\omega(k) = \sqrt{gk + \frac{\sigma}{\rho}k^3} + u_{eff}(k)k$$

Doppler shift

$$\text{With } u_{eff}(k) = 2k \int_0^{\infty} u_L(z) e^{-2kz} dz$$

(Stewart and Joy 1974)



Waves with different **wave-number**
“see” velocities at **different depths**

Non-linear waves' propagation and growth across scales

Part I: Kinematics of modes k

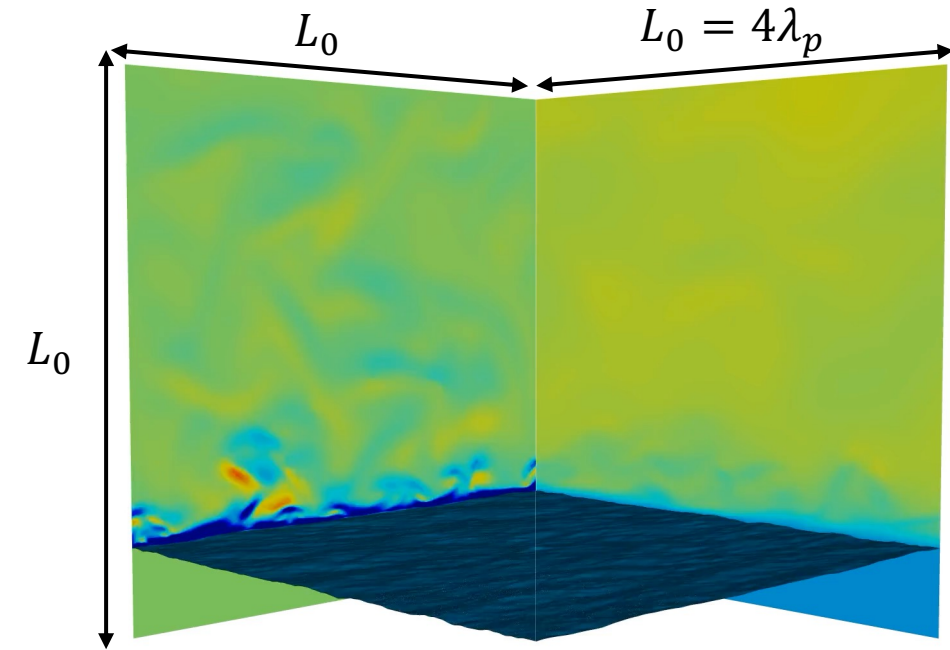
How are non-linear
waves Doppler
shifted?

Part II: Dynamics of modes k

How non-linear
Doppler shifted
waves **grow?**

DNS Two-Phase Flow setup (Adapted from Wu et al. 2022)

Temporally evolving **broadbanded wavefield**, coupled with **a turbulent boundary layers** on both the **air** and the **water**



$t = 0.000$

$u_*/c = 0.25, 0.5$

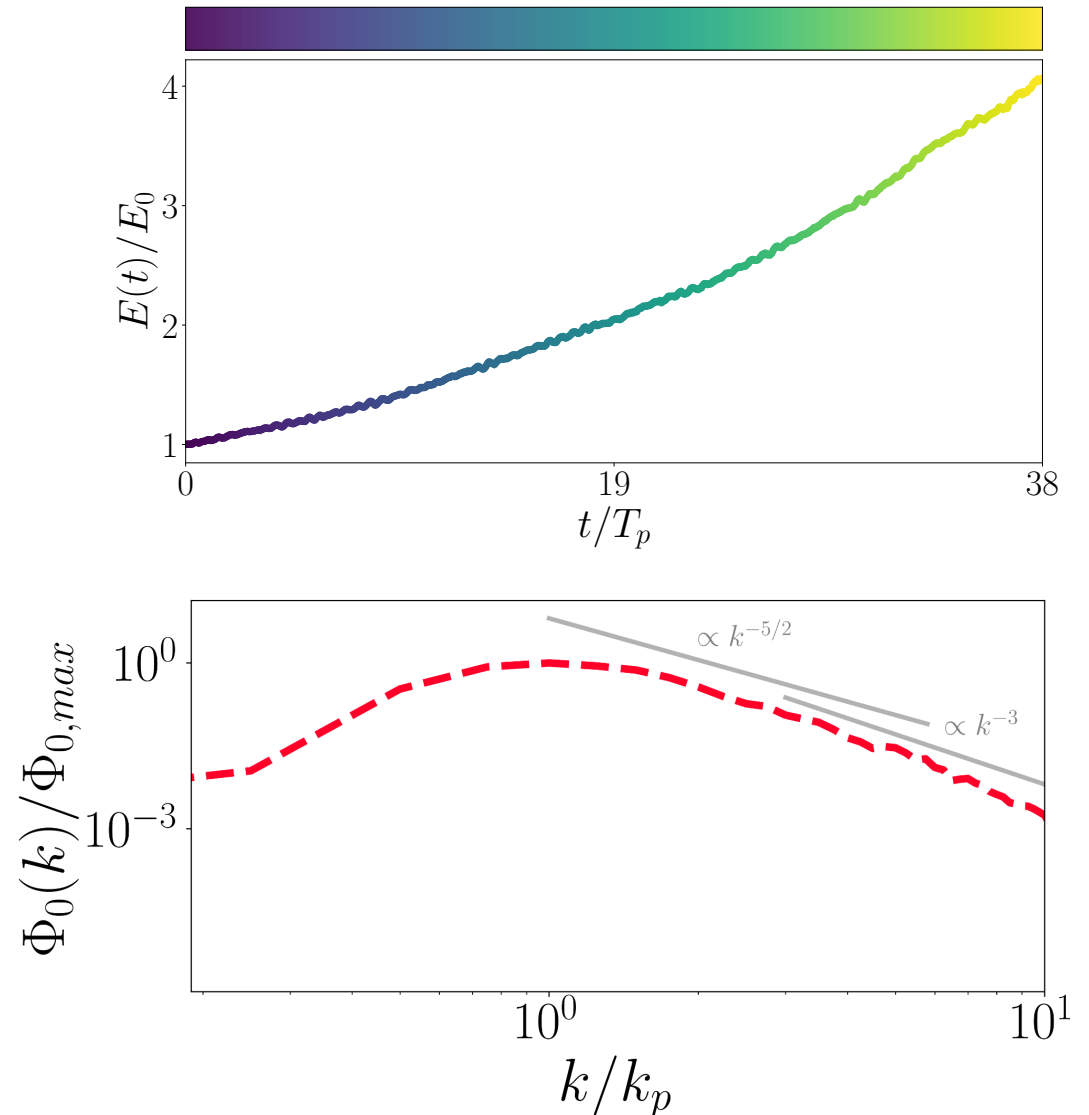
$k_p H_s = 0.08, 0.16$

$$Re_w = \frac{\rho_w c \lambda_p}{\mu_w} = 2.5 \cdot 10^4$$

$$Bo = \frac{(\rho_w - \rho_a)g}{k_p^2 \sigma} = 200$$



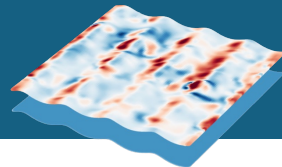
Basilisk: Open-source solver (developed by Stephane Popinet): <http://basilisk.fr/>



Non-linear waves' propagation and growth across scales

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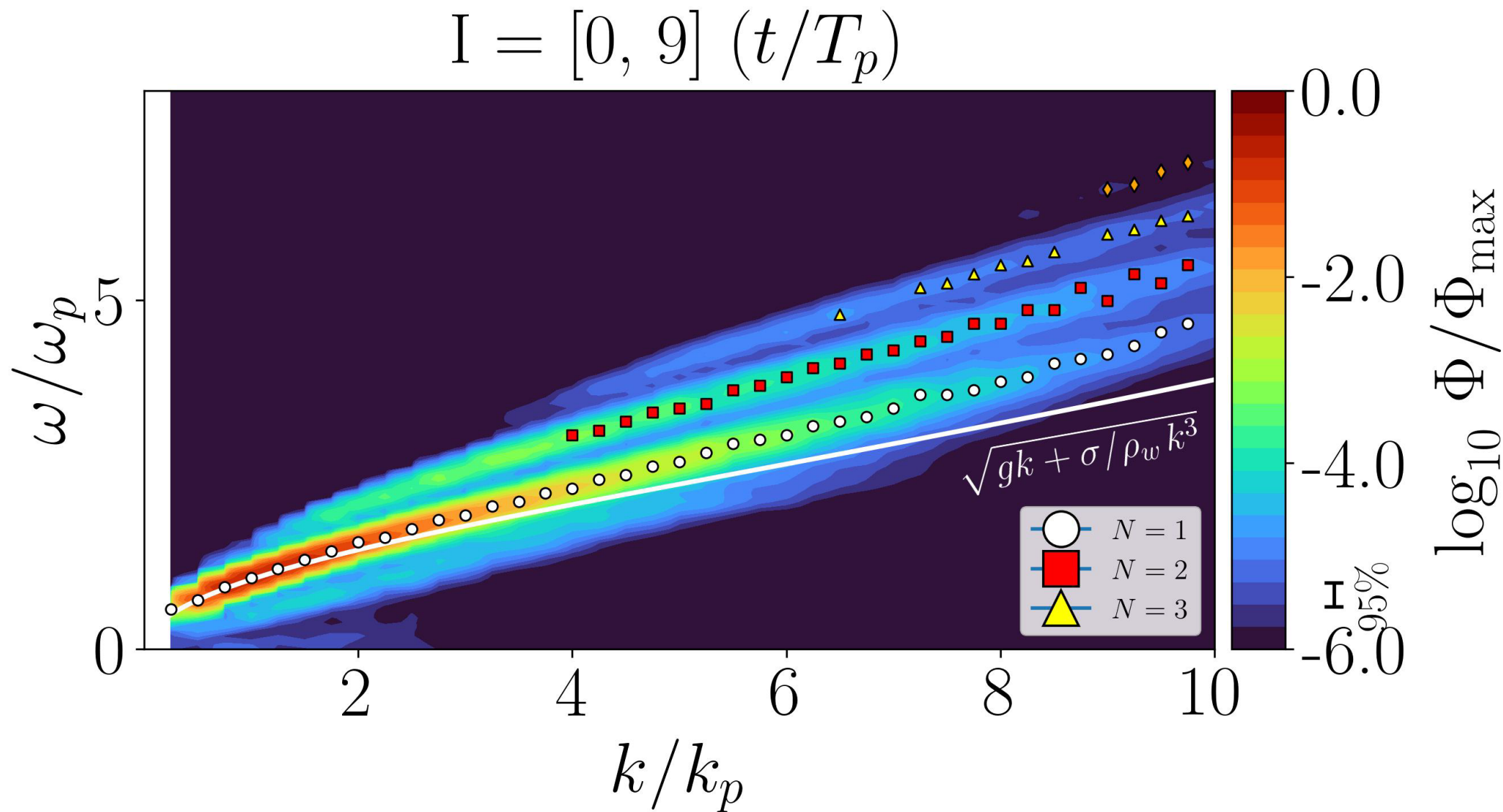
How are non-linear
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$$\eta(x, y, t)$$

FFT

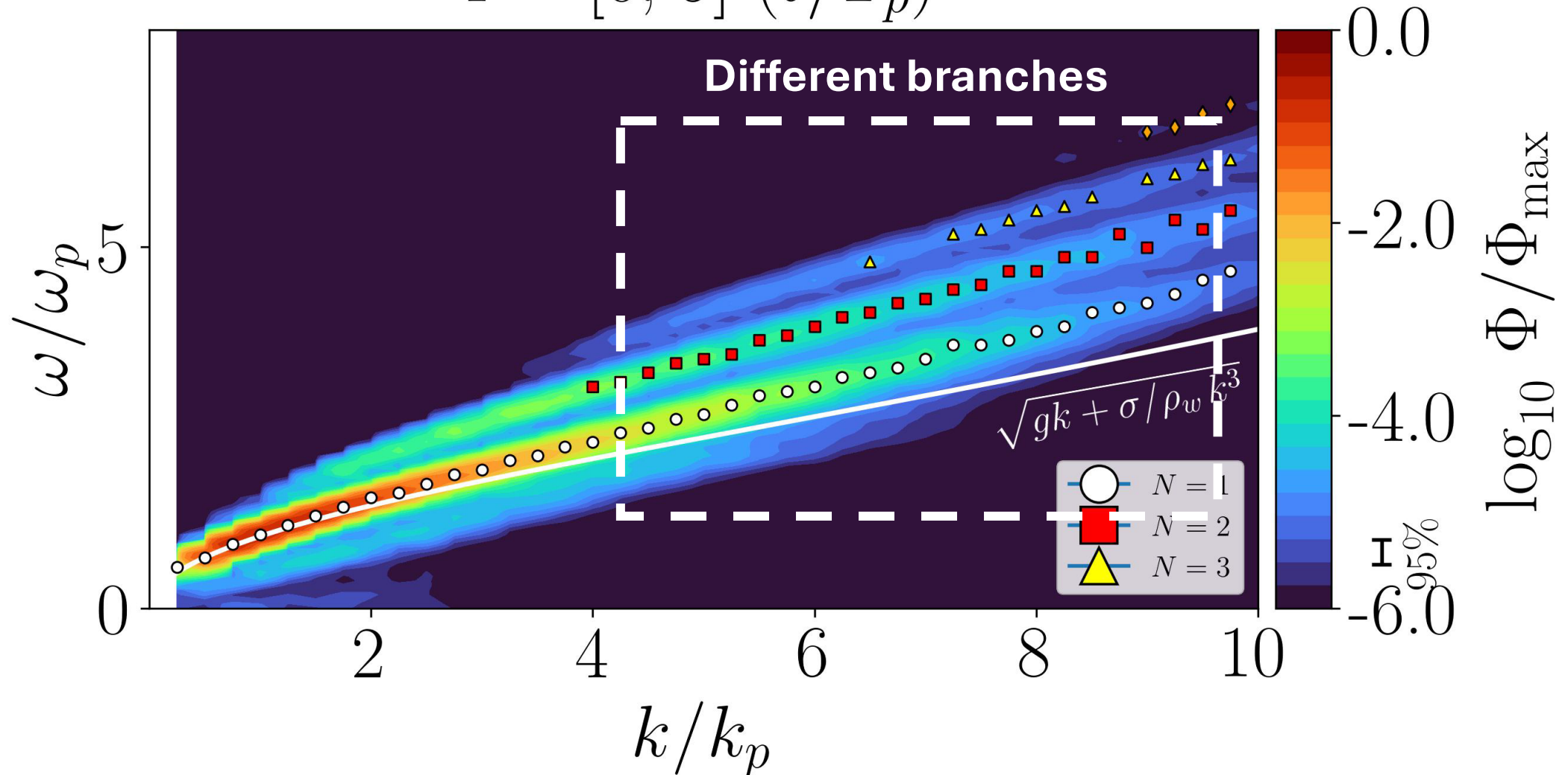
$$\Phi(k, \omega, t)$$



Higher order harmonics

Energy is contained in multiple branches including the linear **dispersion relation**

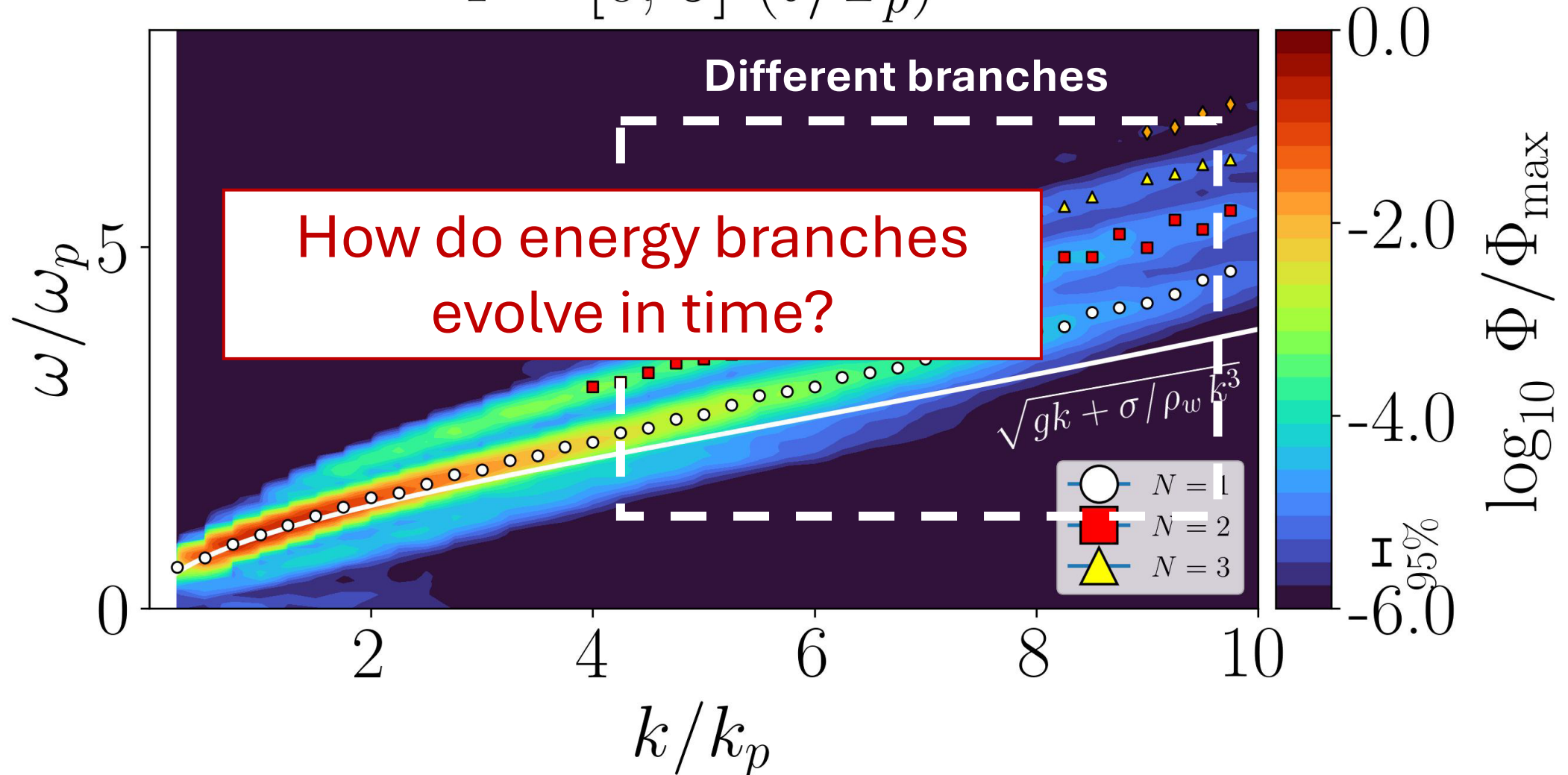
$$I = [0, 9] \quad (t/T_p)$$



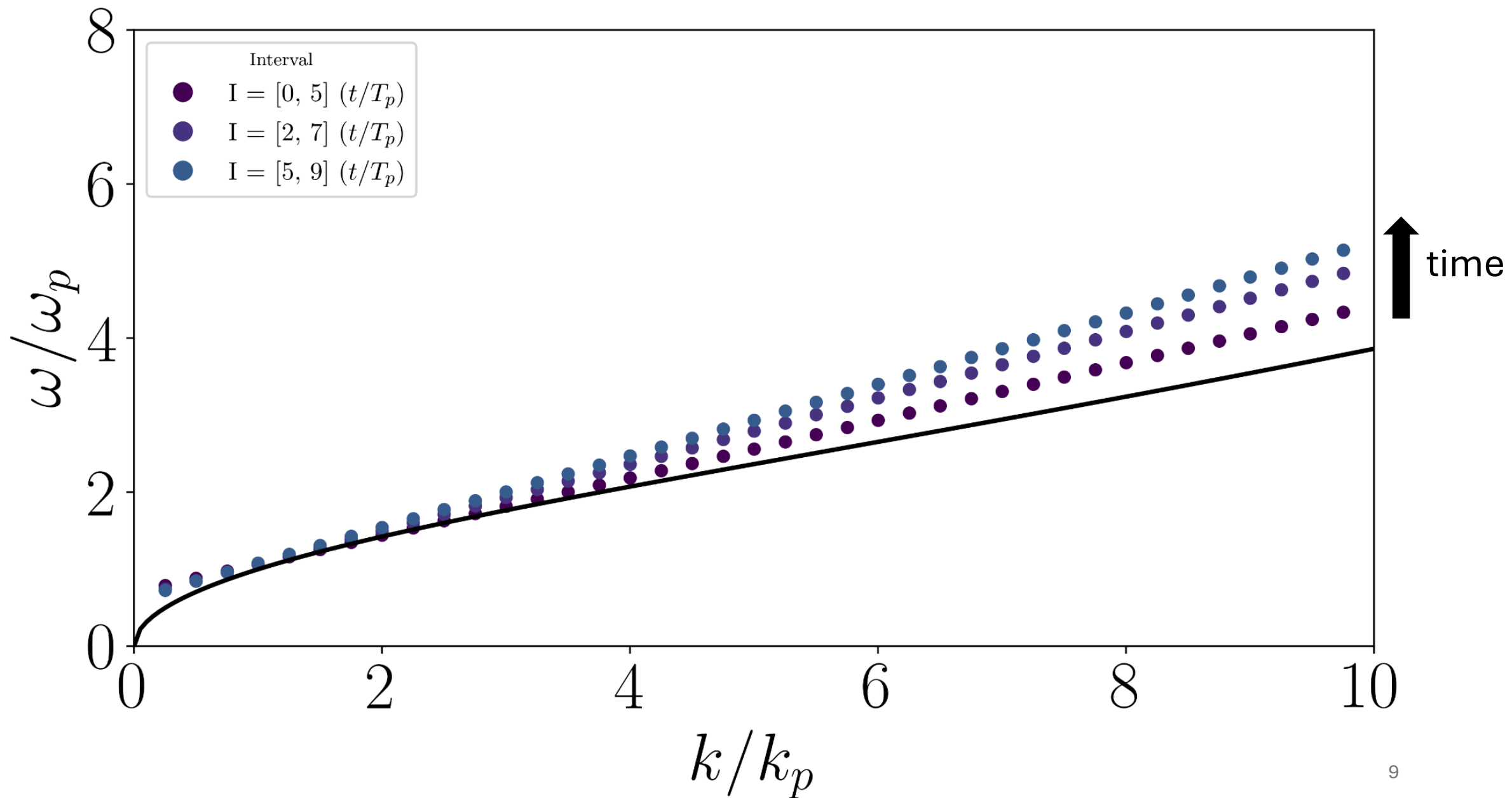
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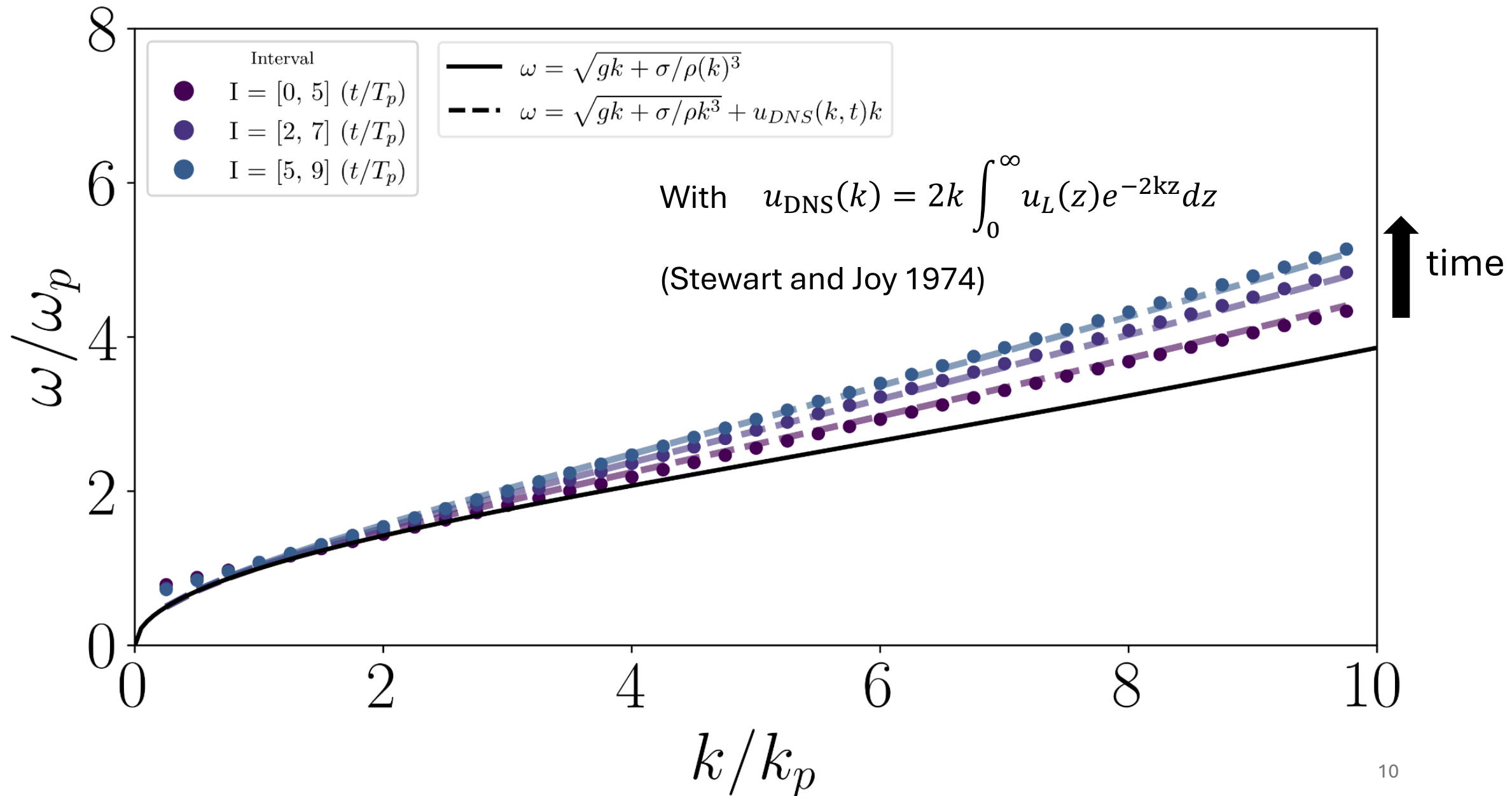
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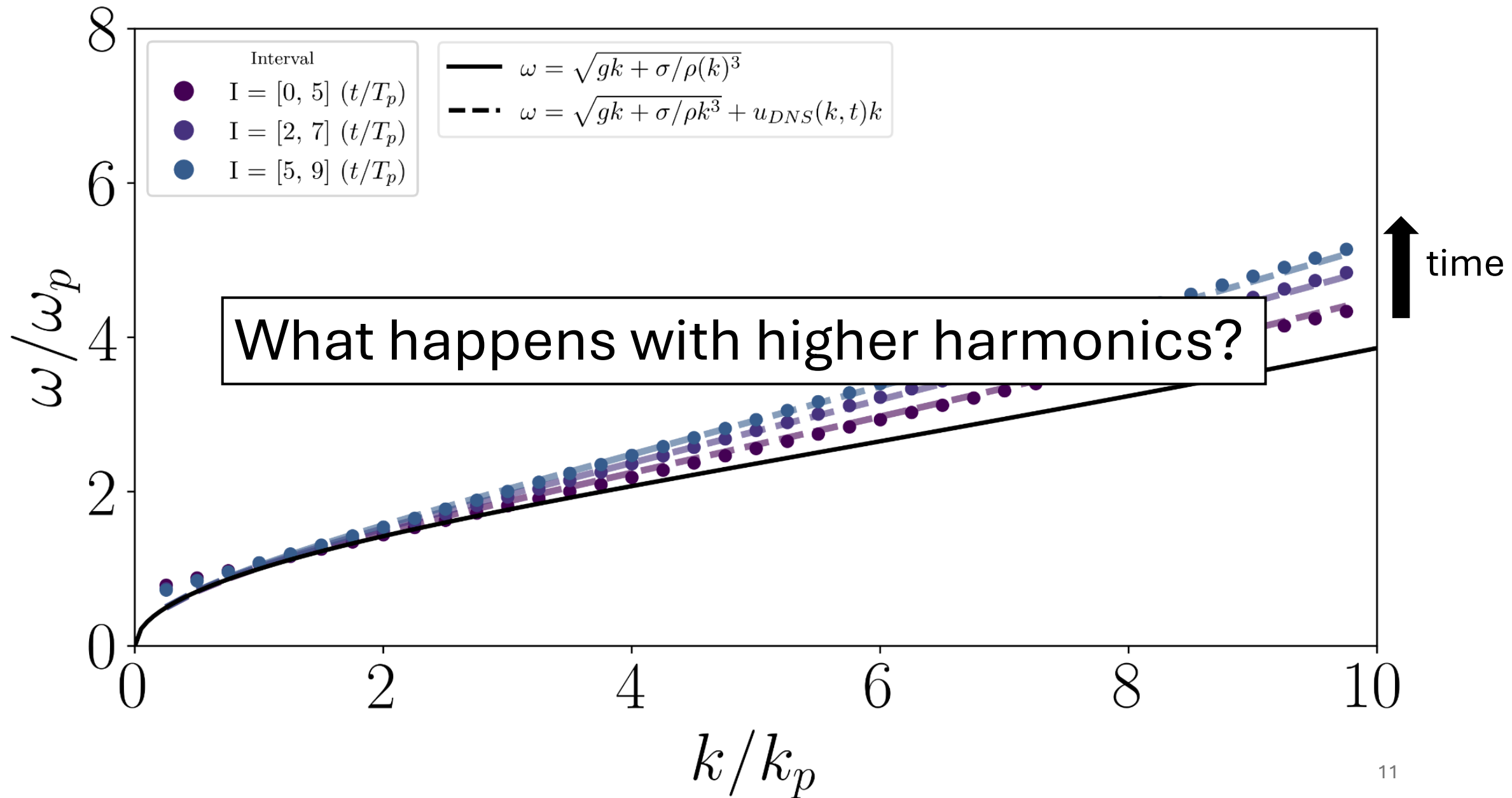
Shift from Linear Dispersion Relationship



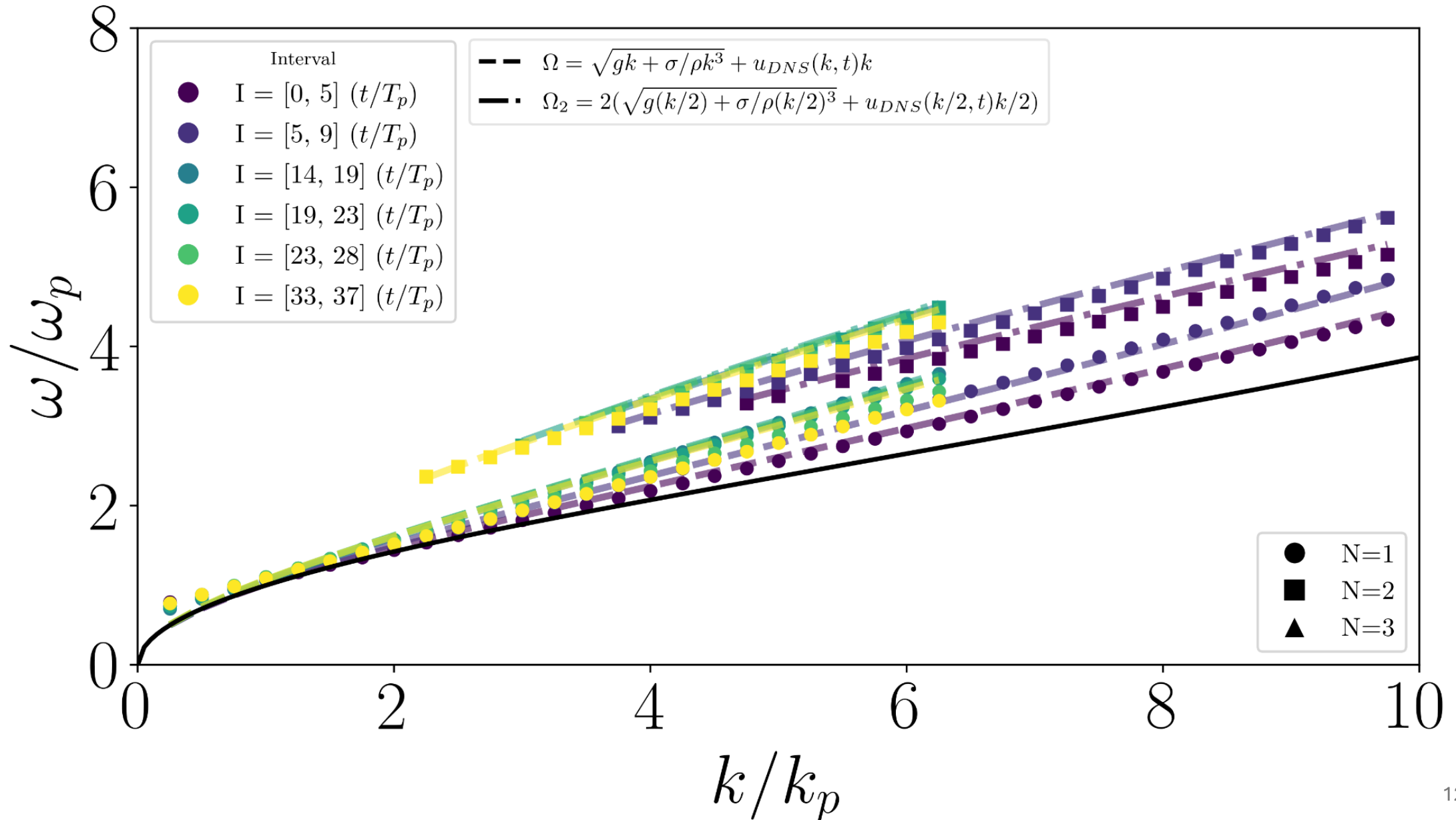
Doppler shift of the first branch from linear dispersion relation



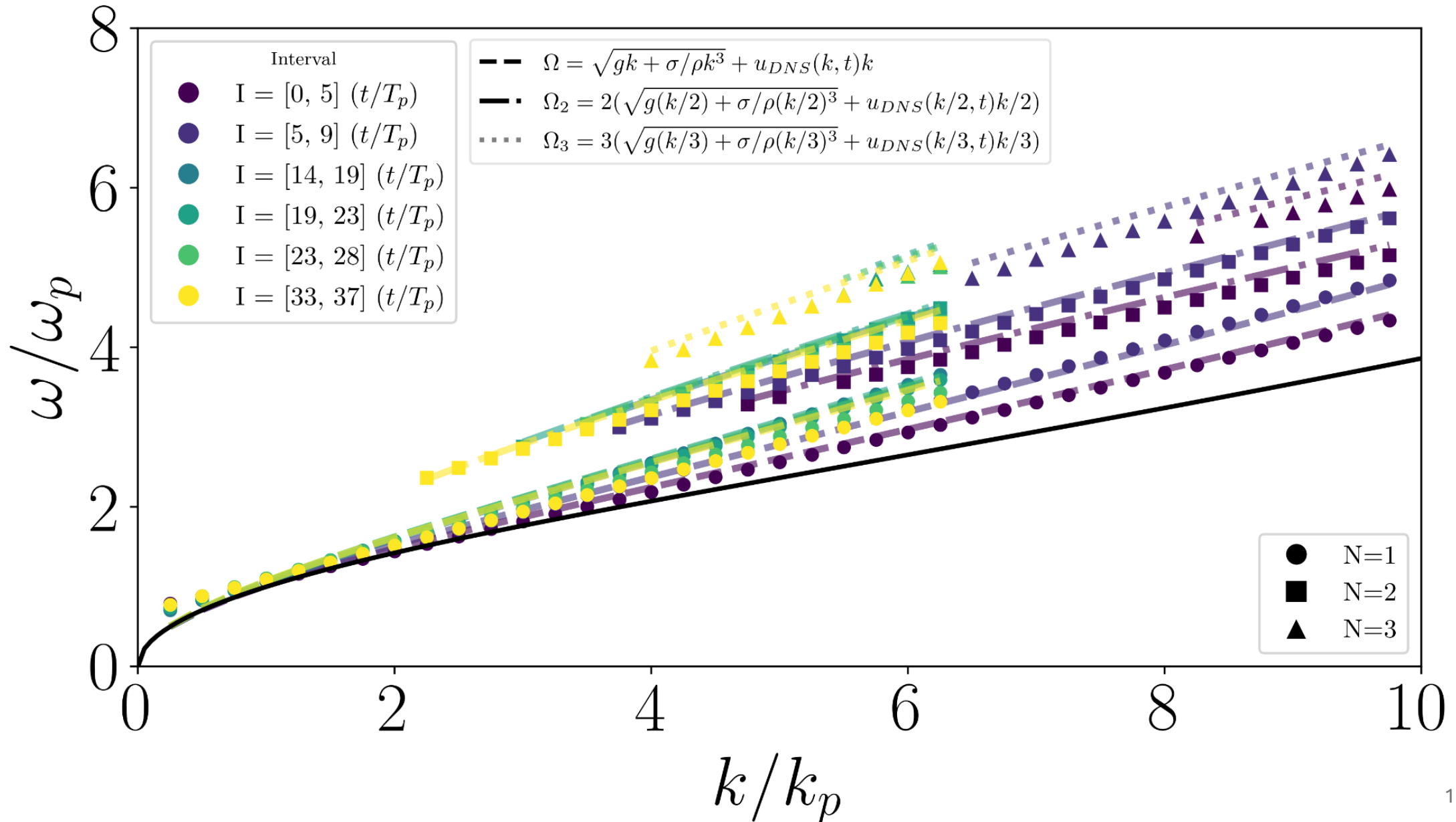
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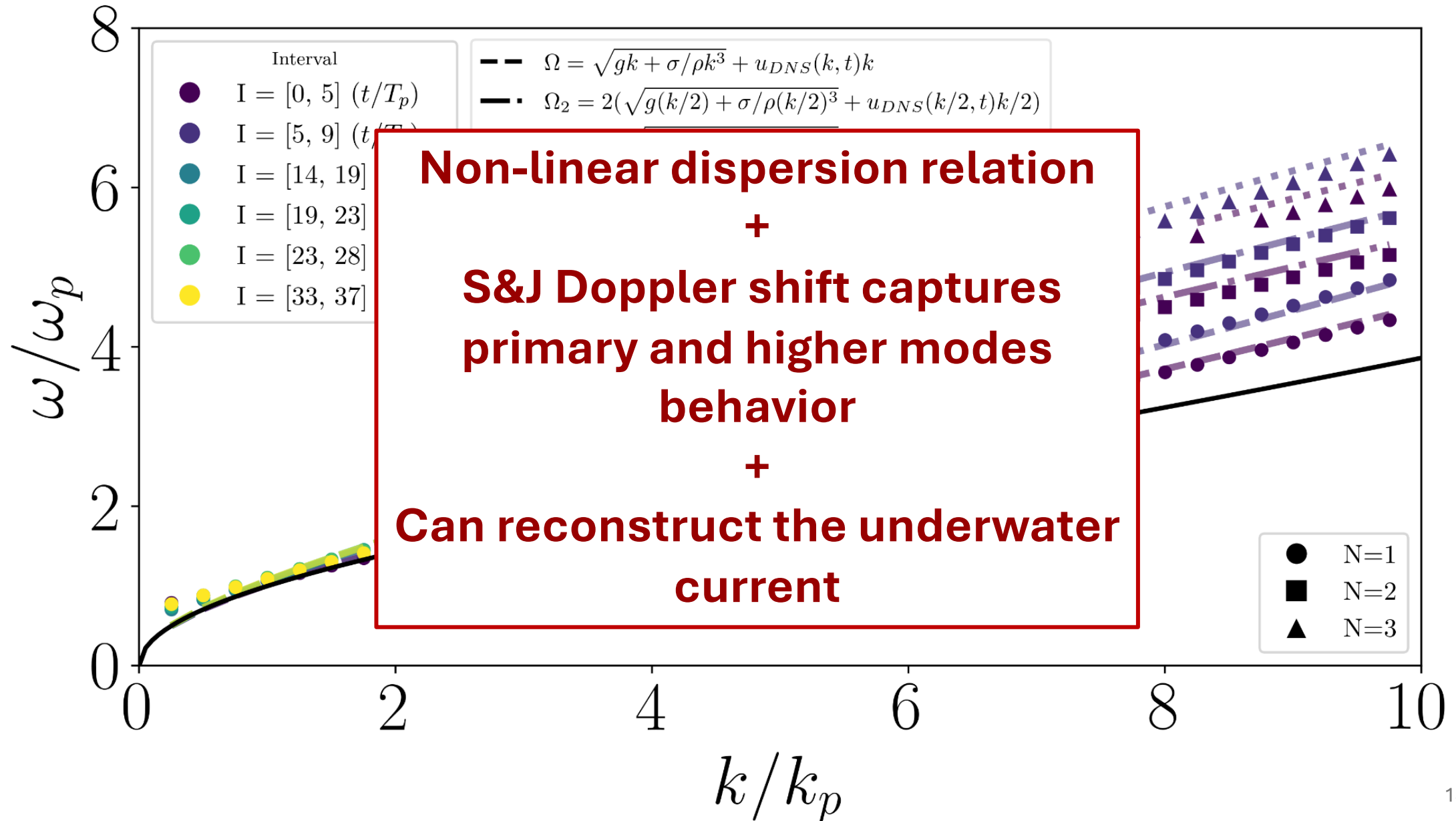
Doppler shift of the other branches following the non-linear dispersion relation



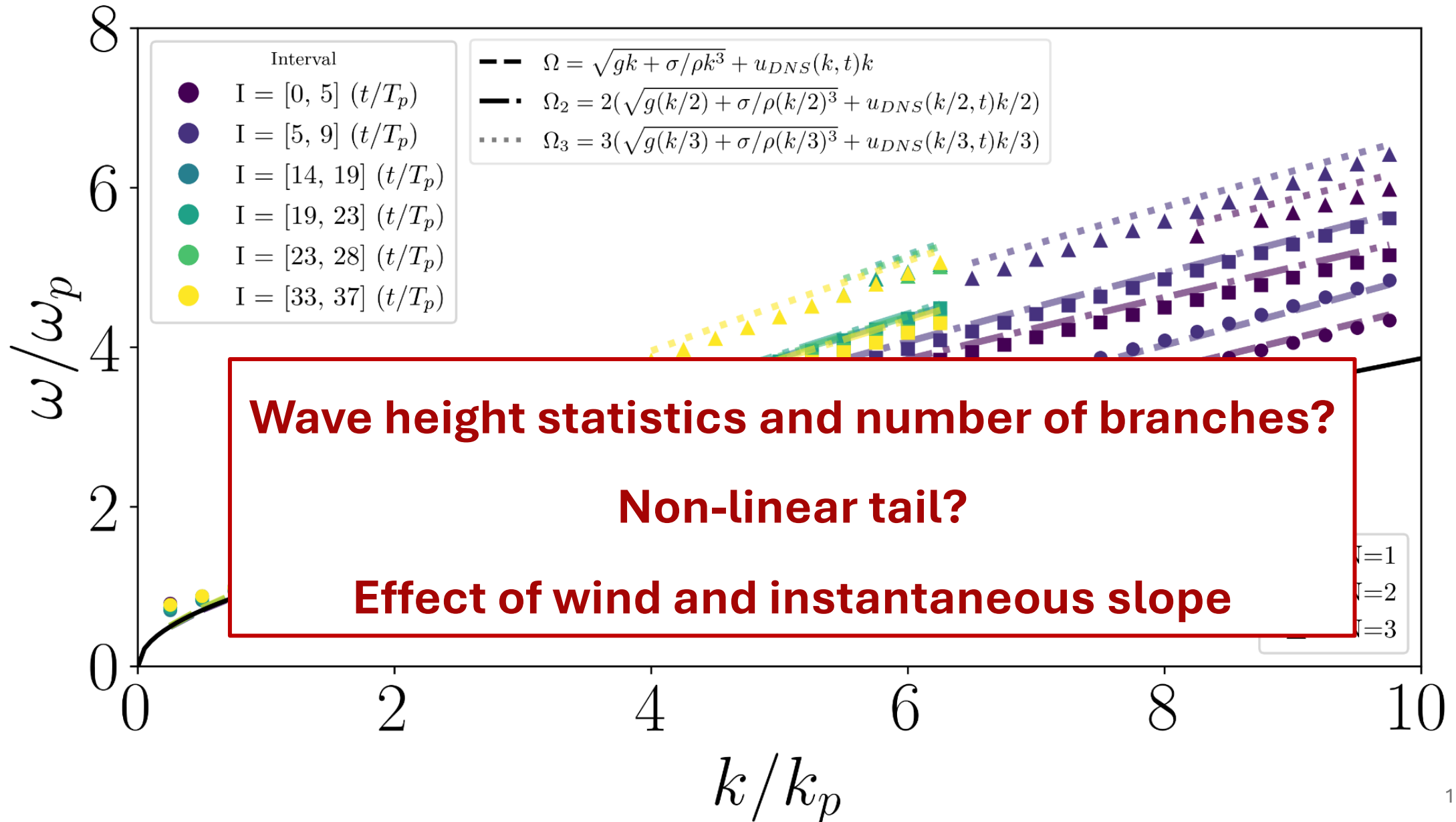
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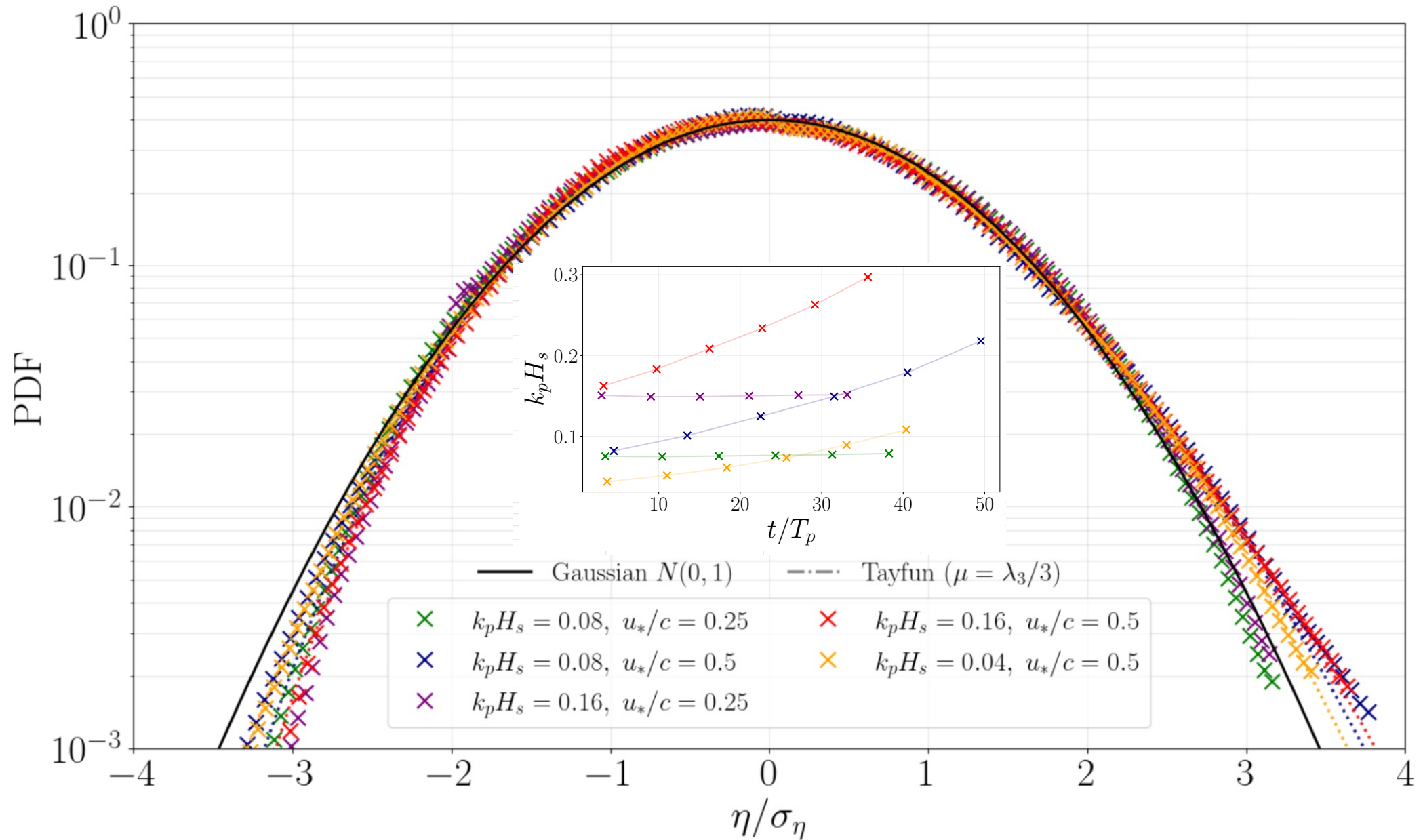
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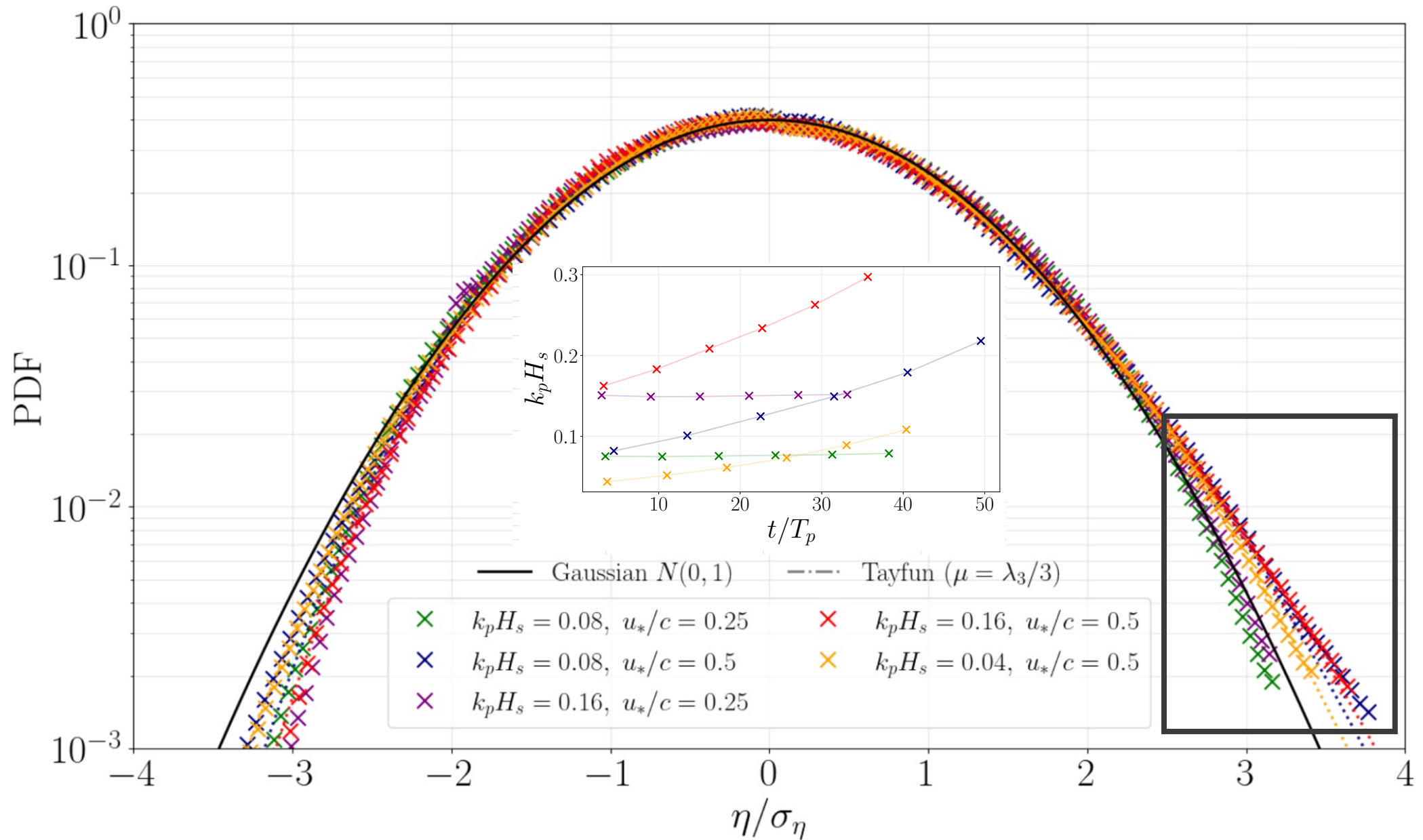
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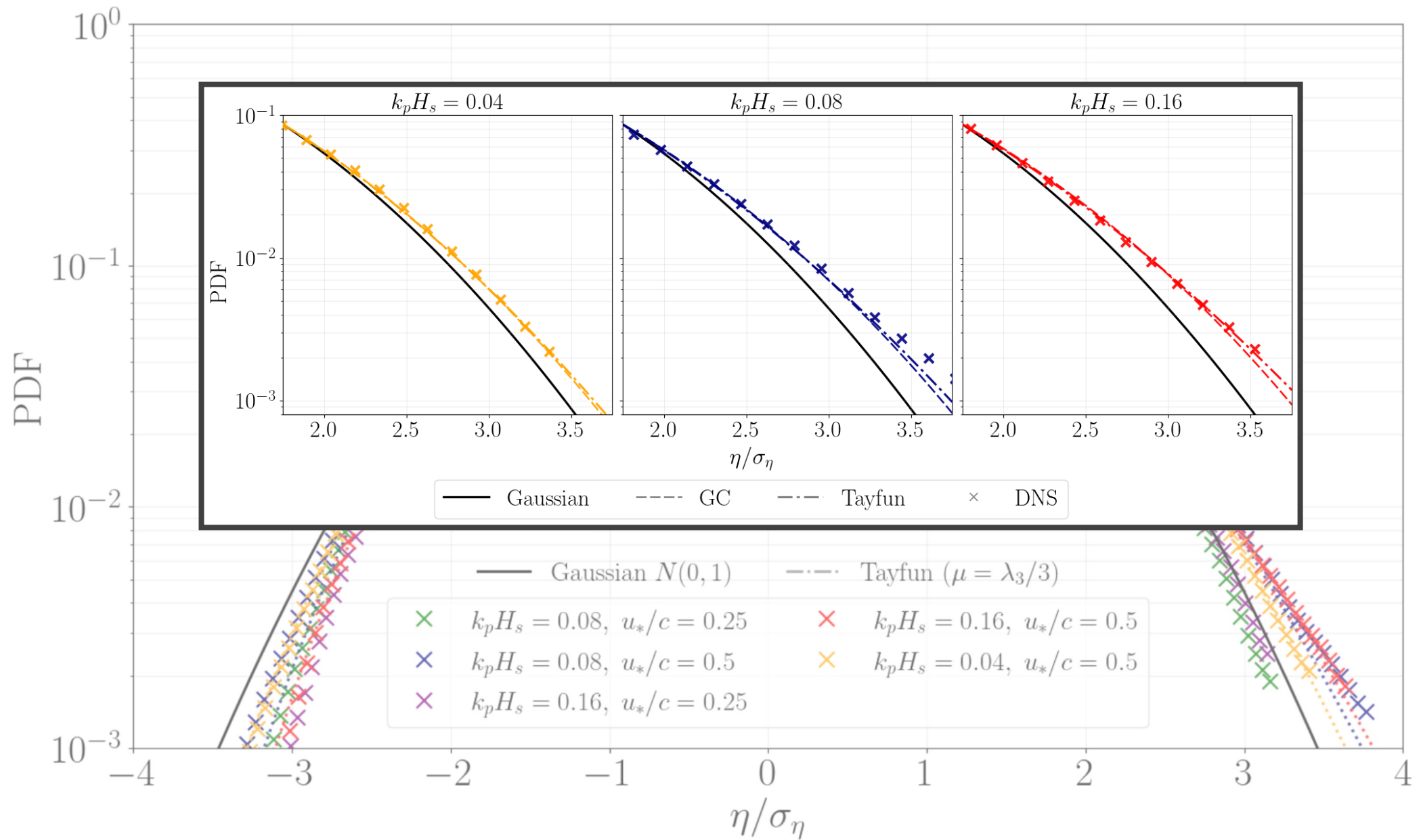
Wave heights statistics



Wave heights statistics



Wave heights statistics



Non-linear waves' propagation and growth across scales



Kinematics of gravity capillary waves above an evolving underwater current

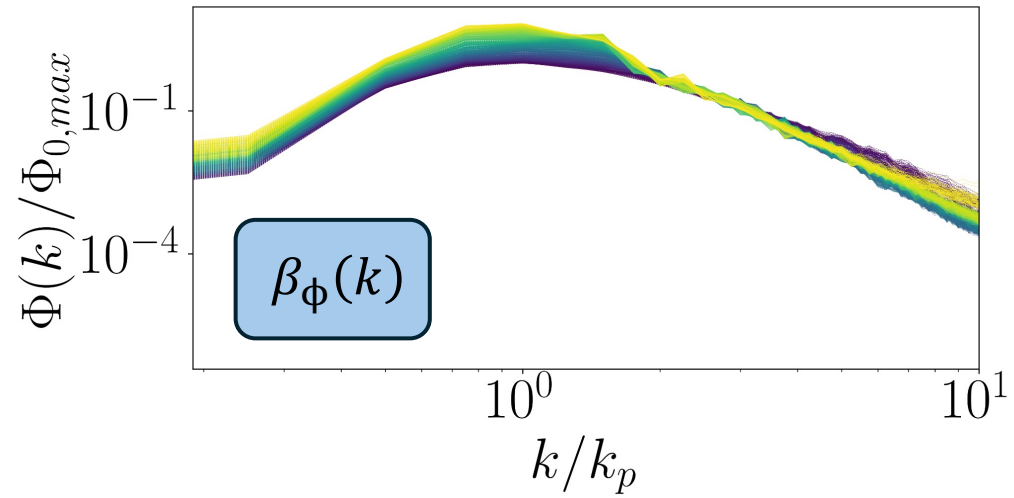
C. Martin Blanco, N. Scapin, J. Wu, S. Popinet,
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Journal of Fluid Mechanics.



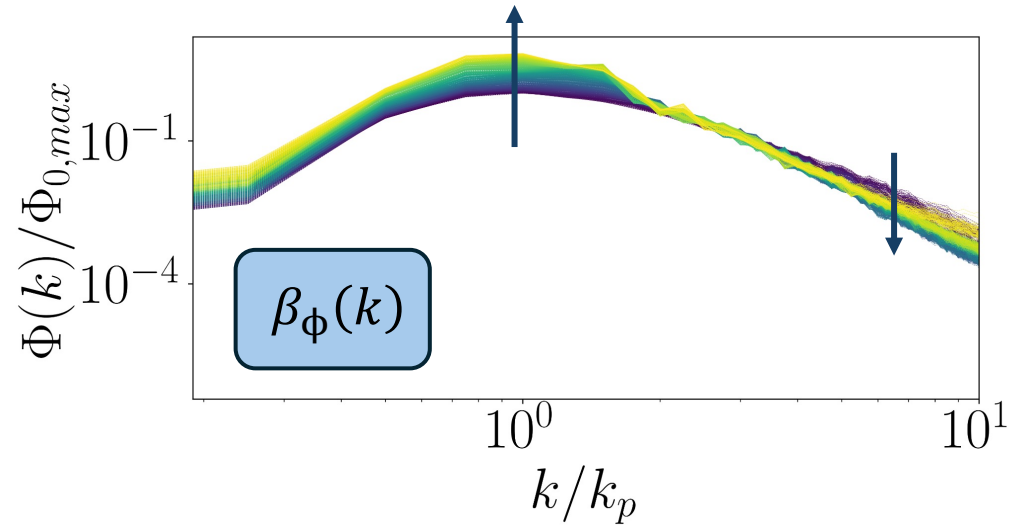
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of modes k

How non-linear
Doppler shifted
waves **grow**?

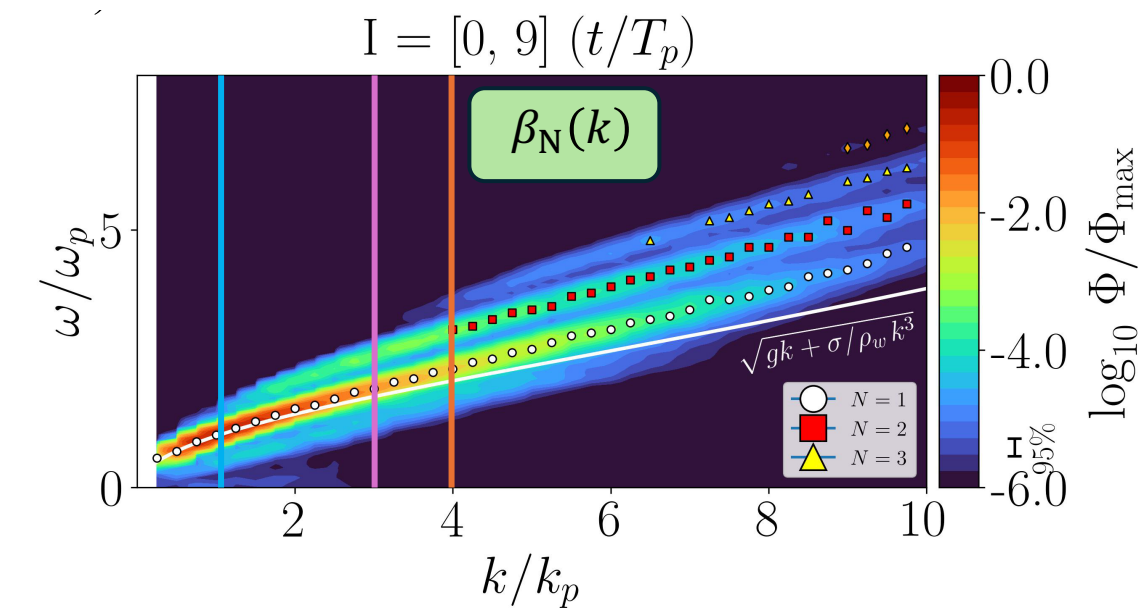
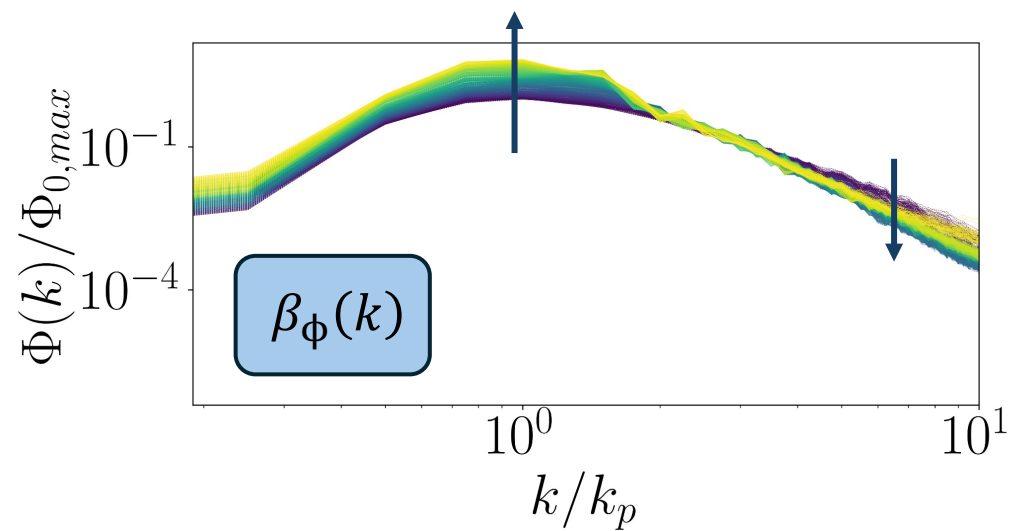
$(k - \omega)$ spectral analysis: growth of modes $\beta_\Phi(k)$ and $\beta_N(k)$



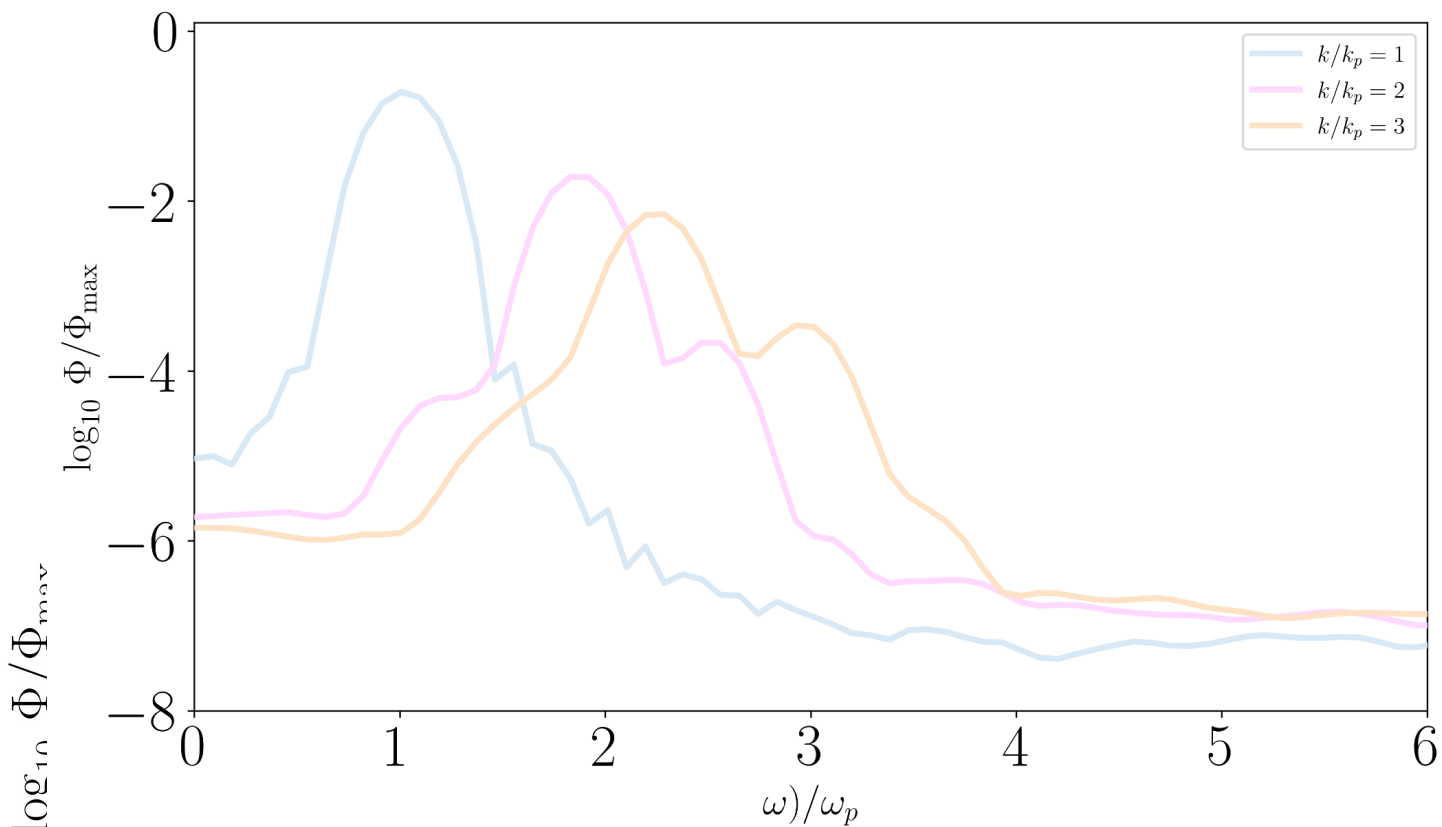
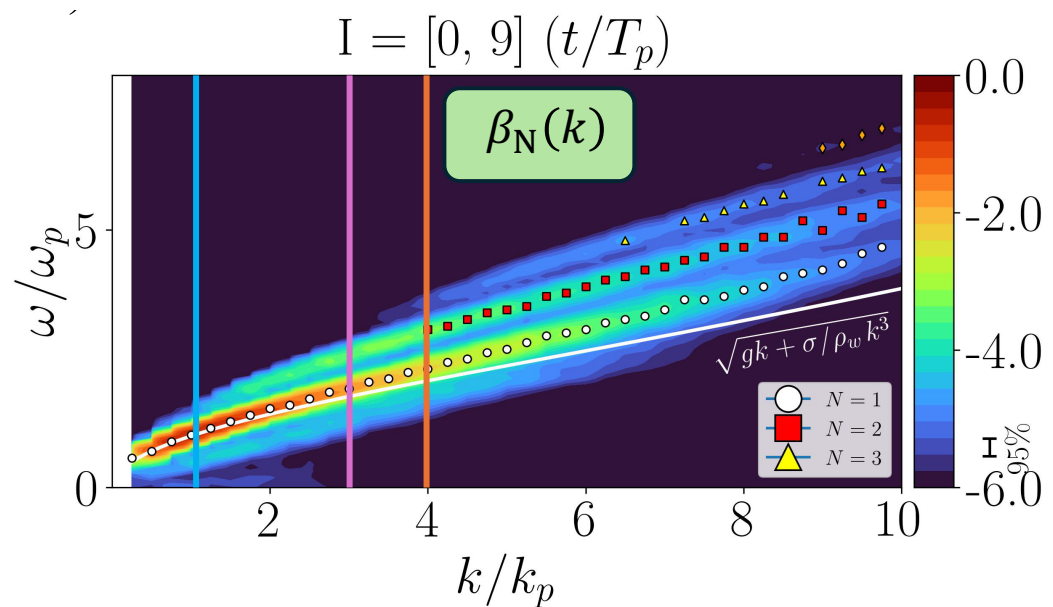
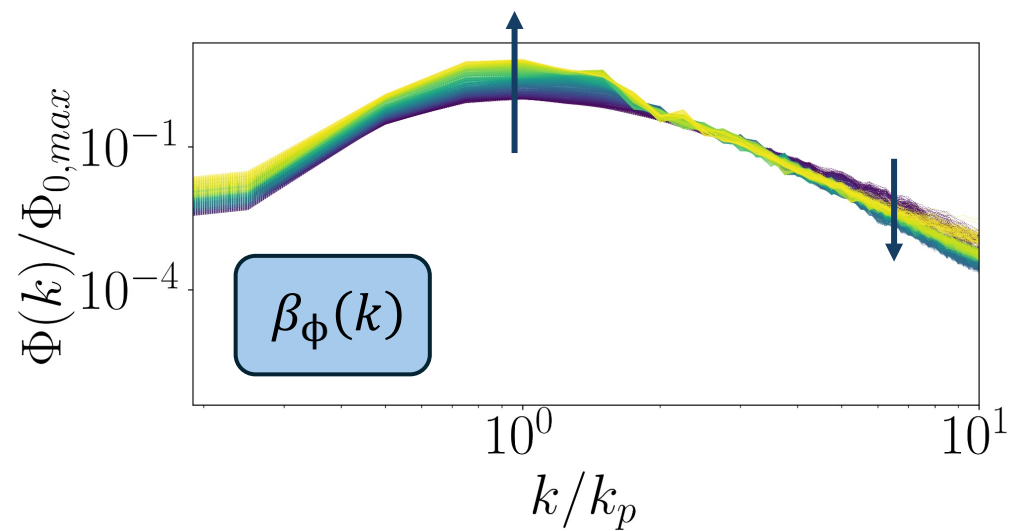
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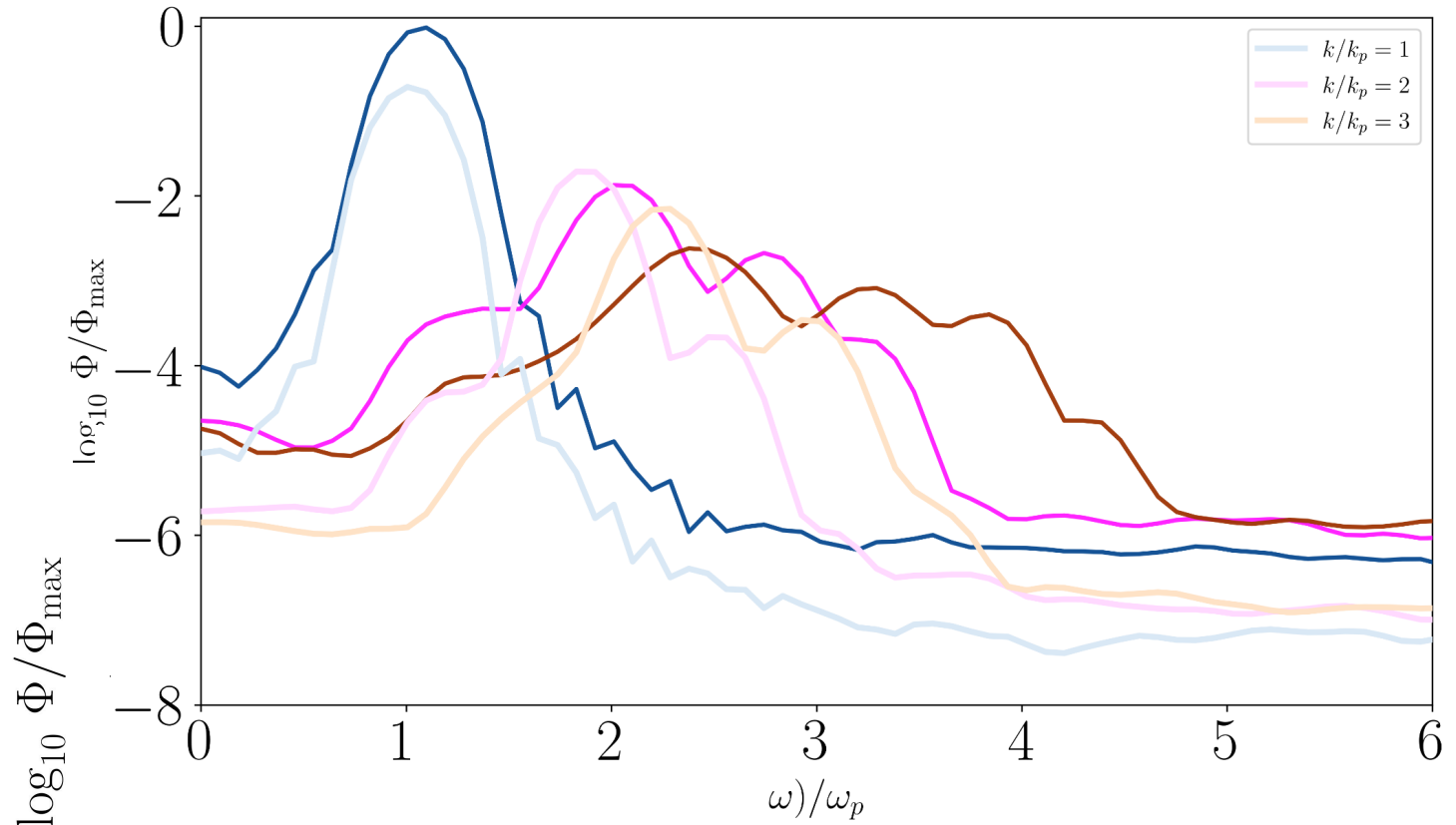
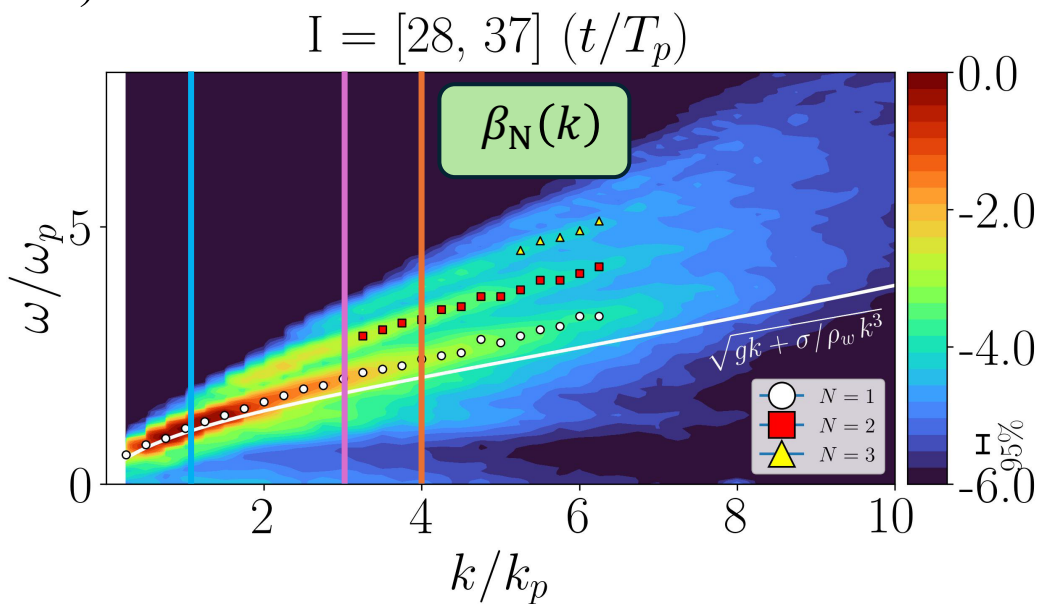
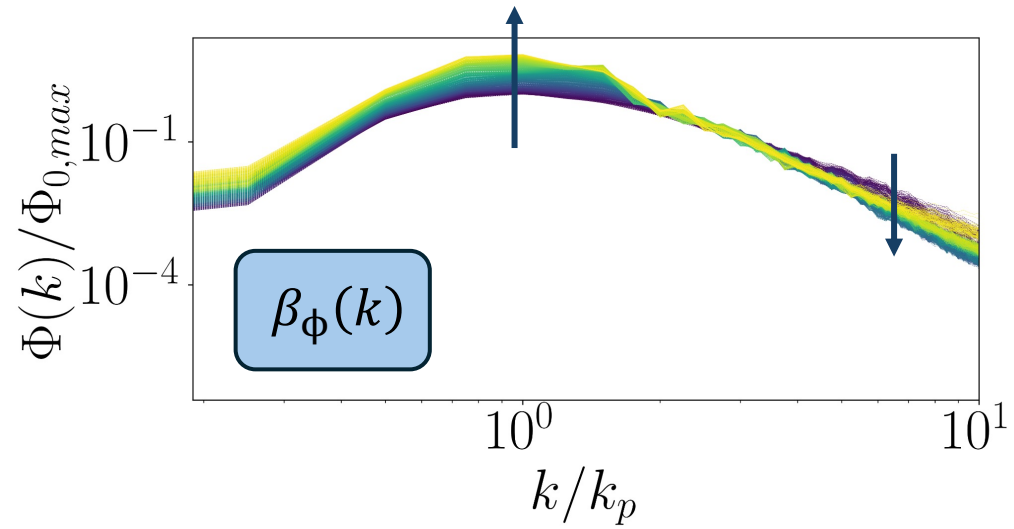
($k - \omega$) spectral analysis: total growth rates $\beta_\Phi(k)$ and growth rate per branch $\beta_N(k)$



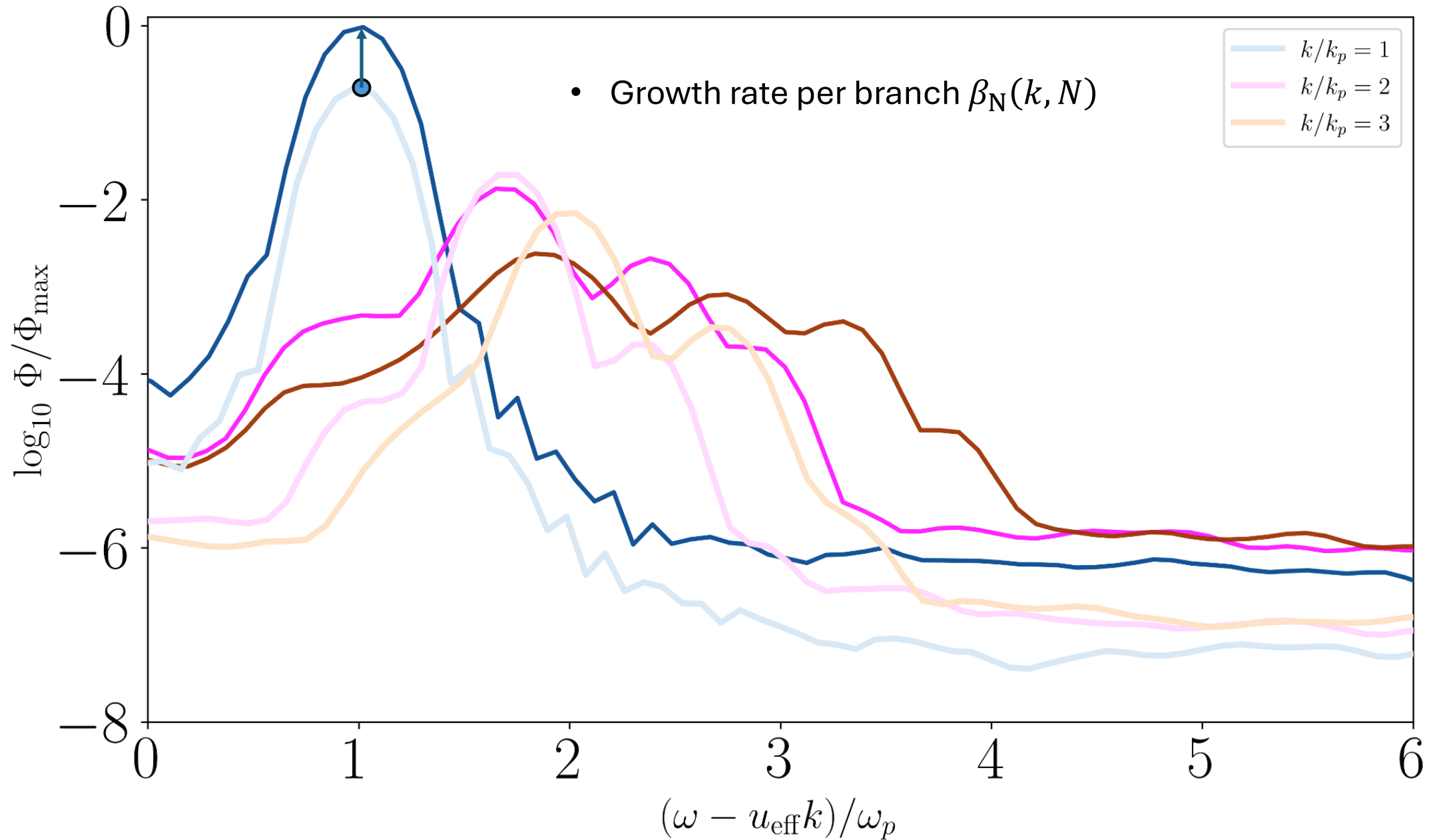
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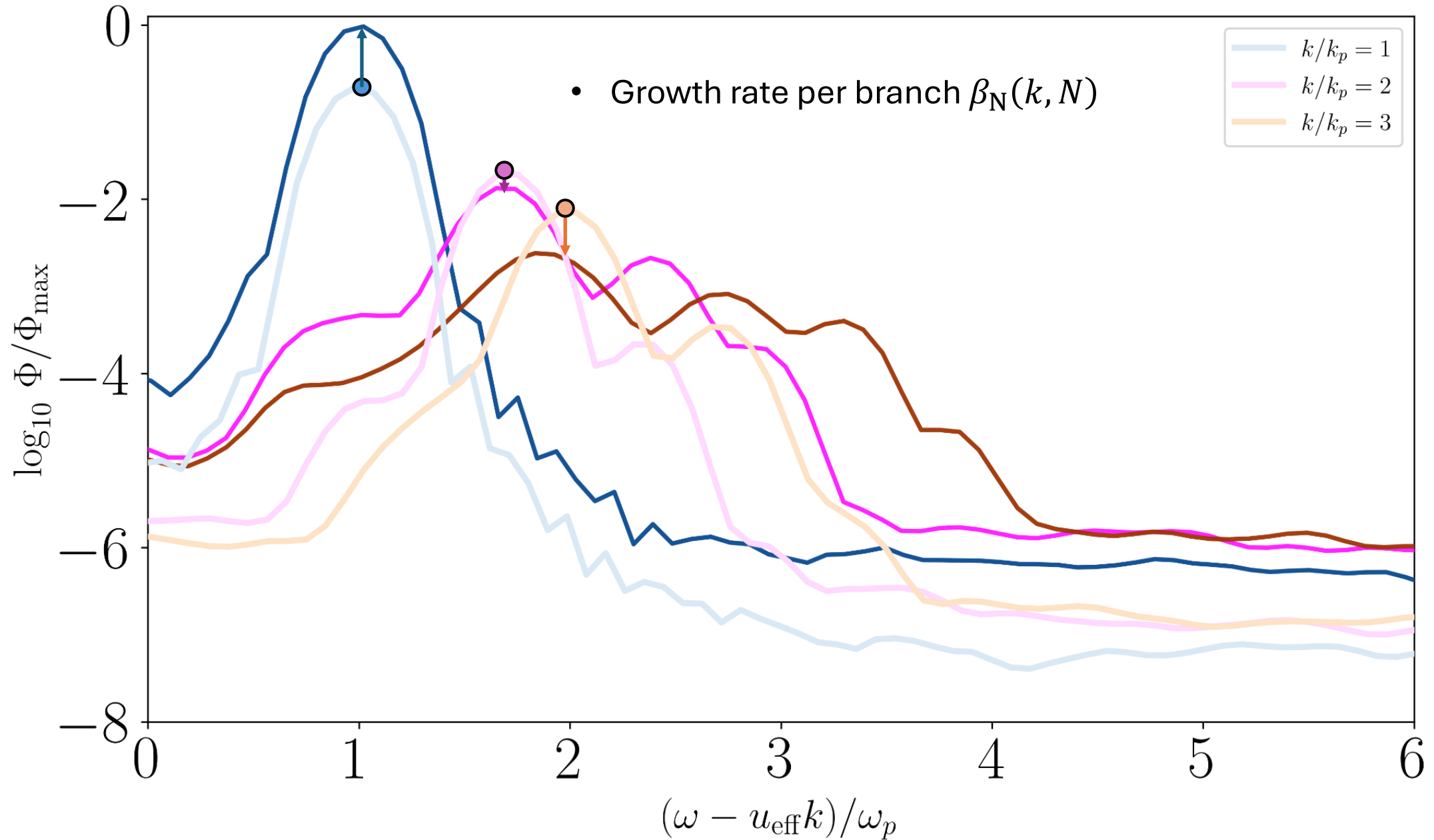
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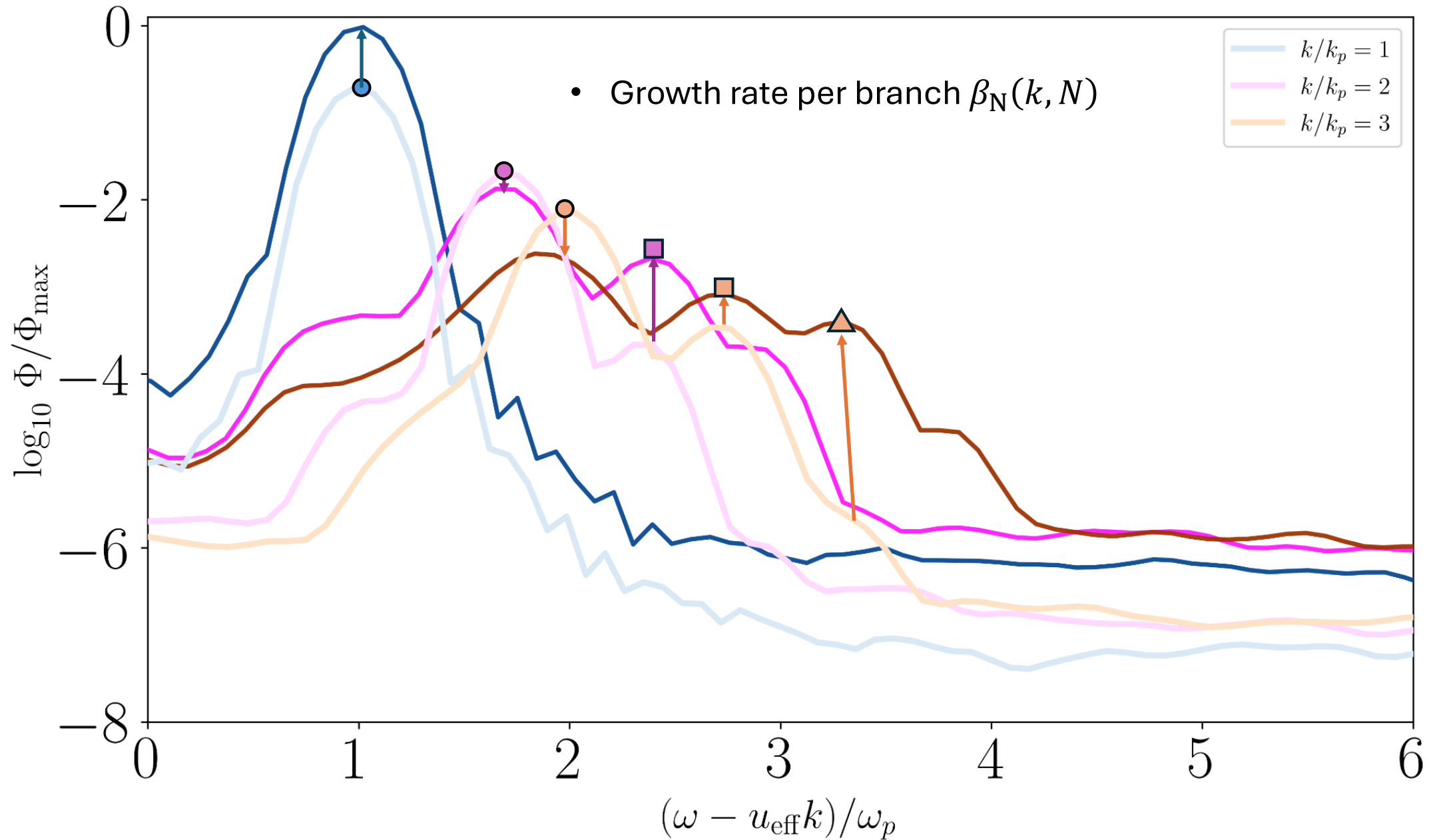
Energy Extraction: Growth rate per branch accounting for Doppler shift $\beta_N(k, N)$



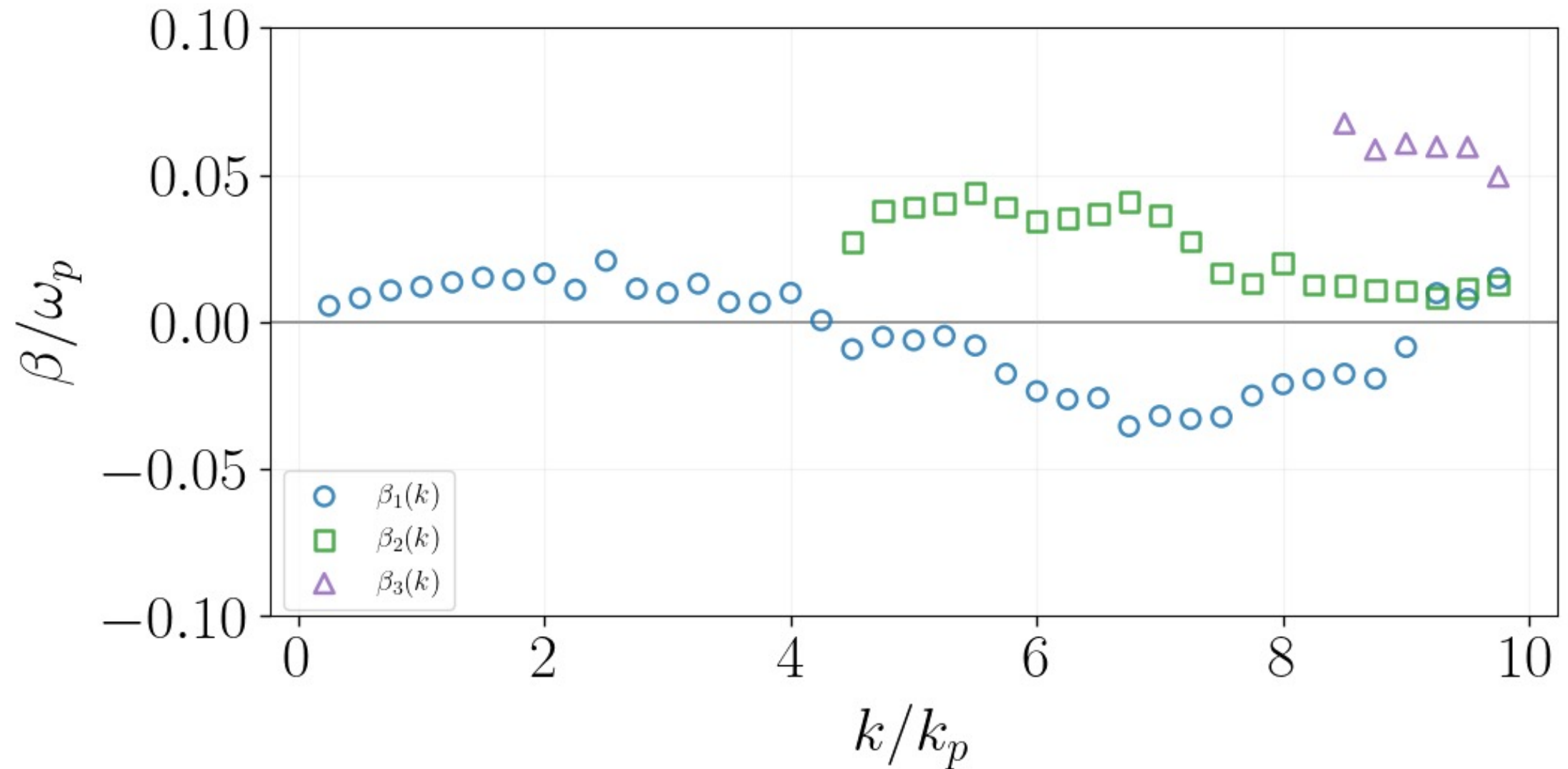
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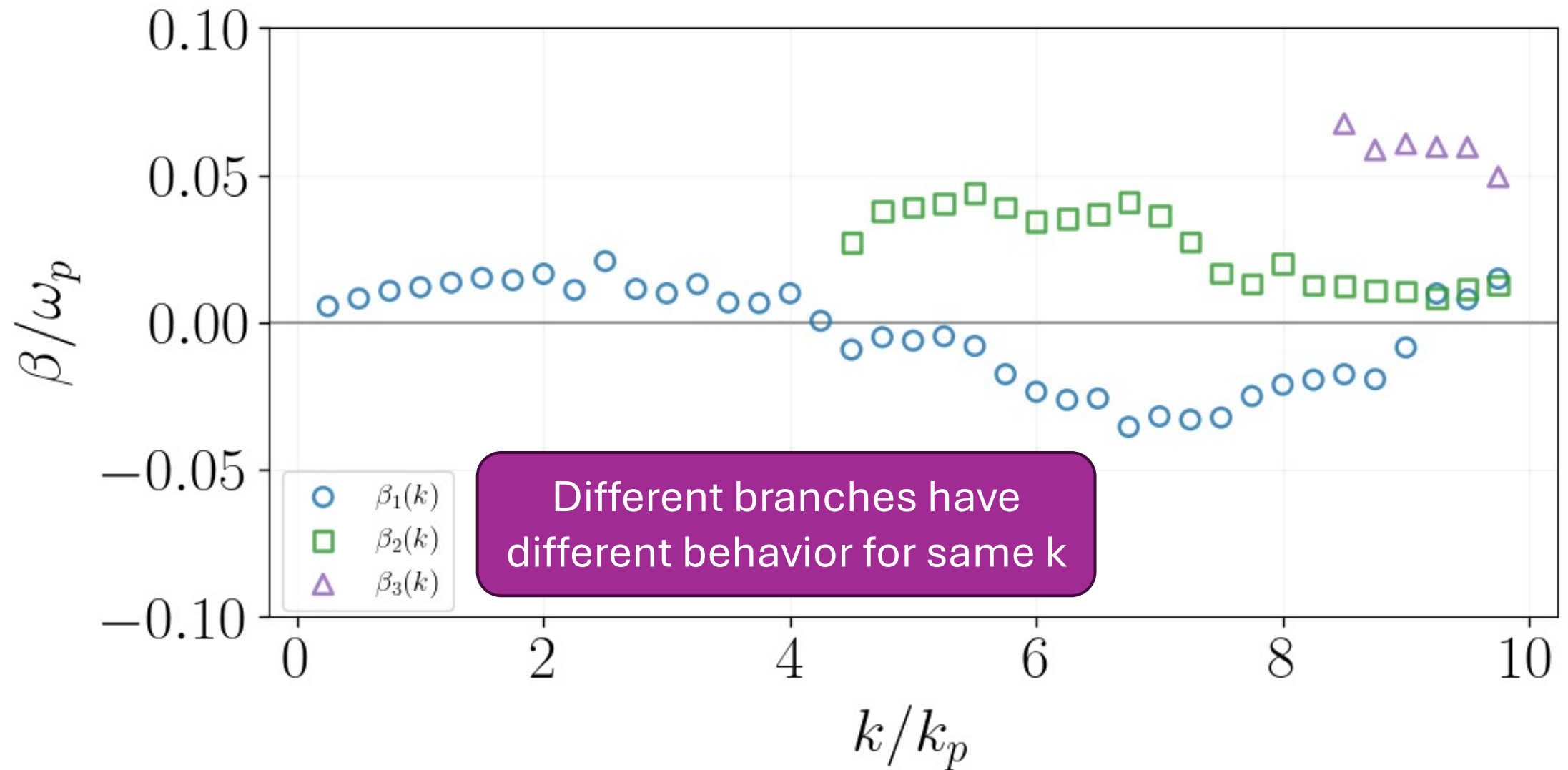
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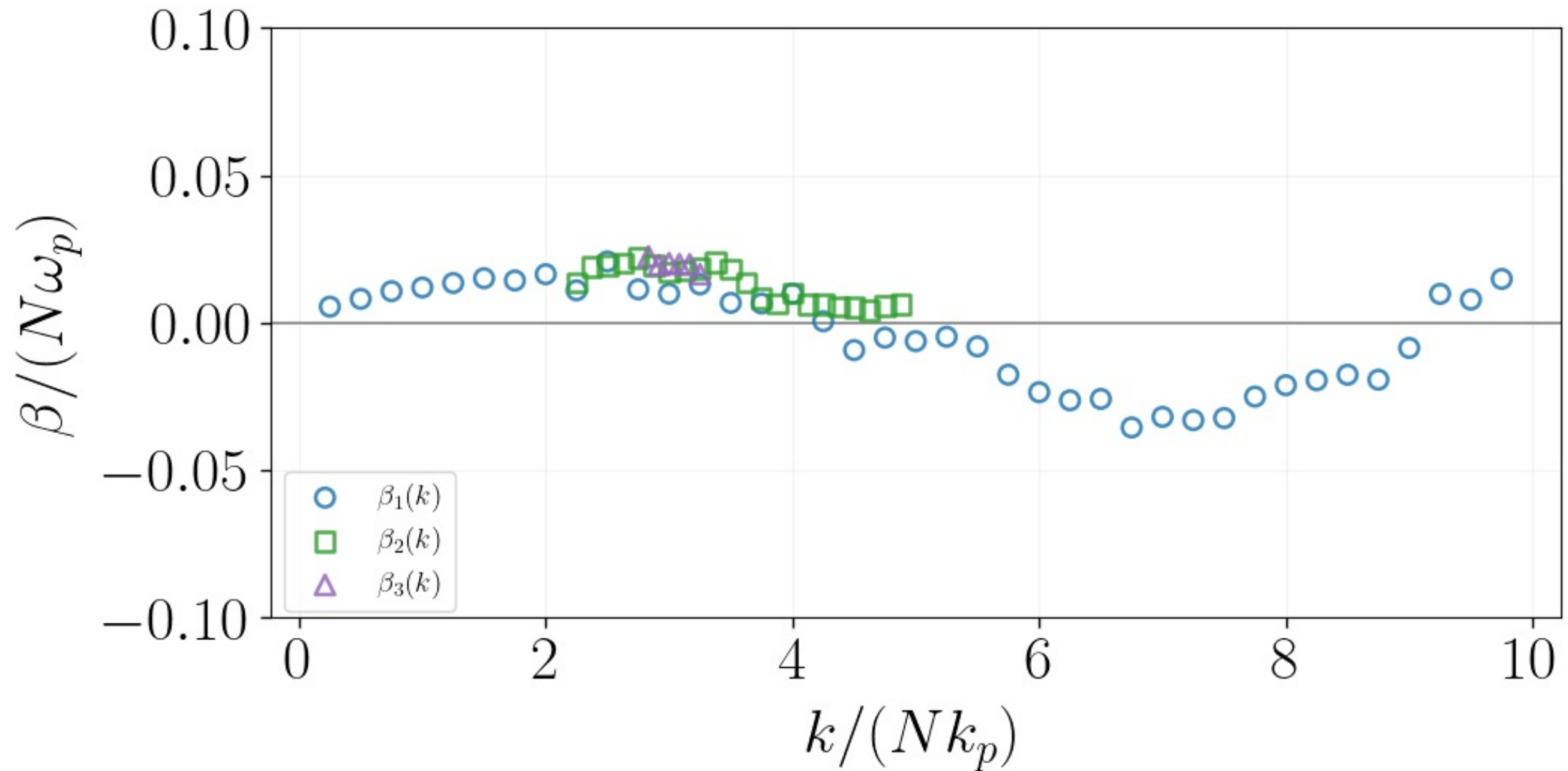
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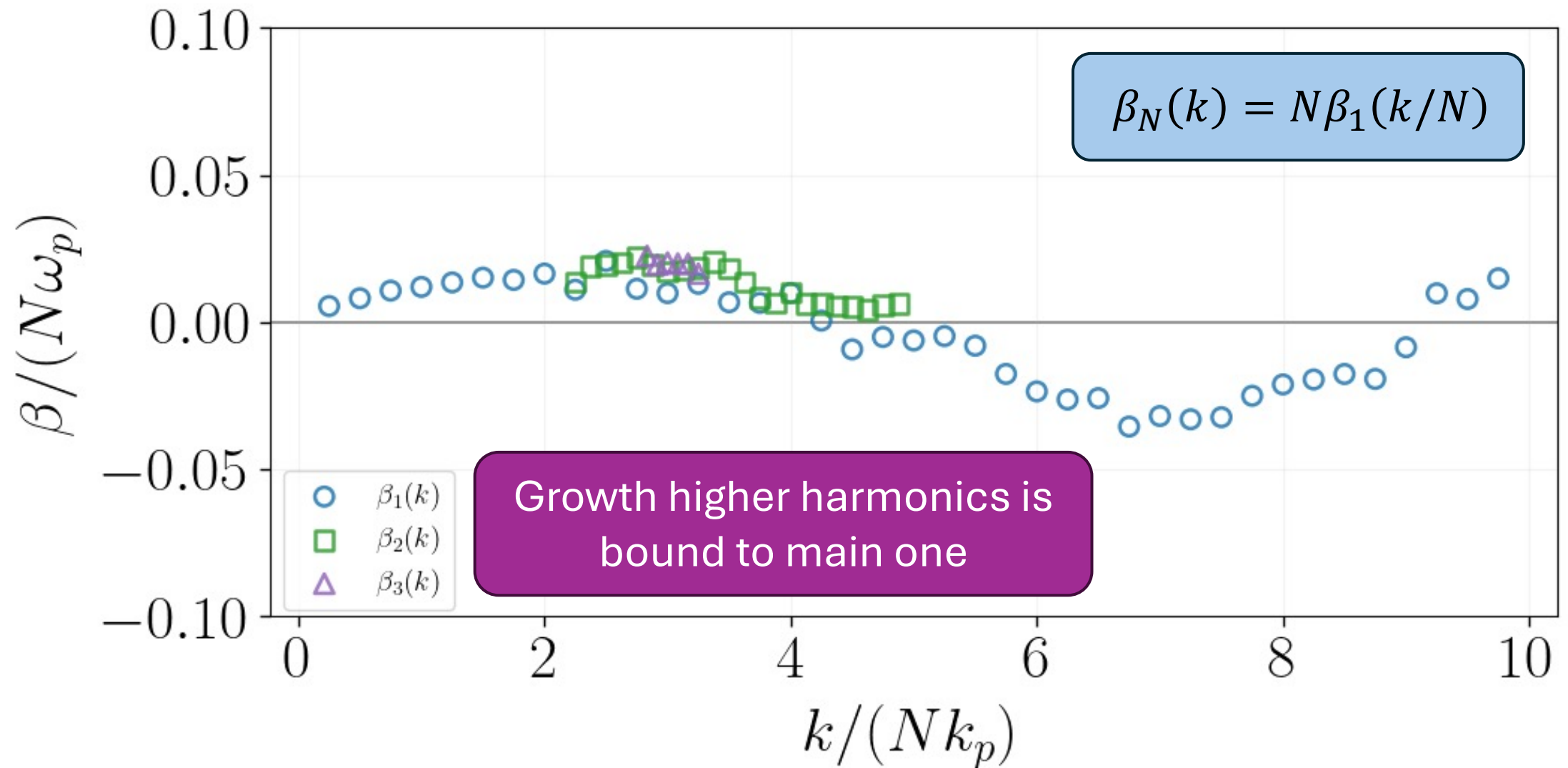
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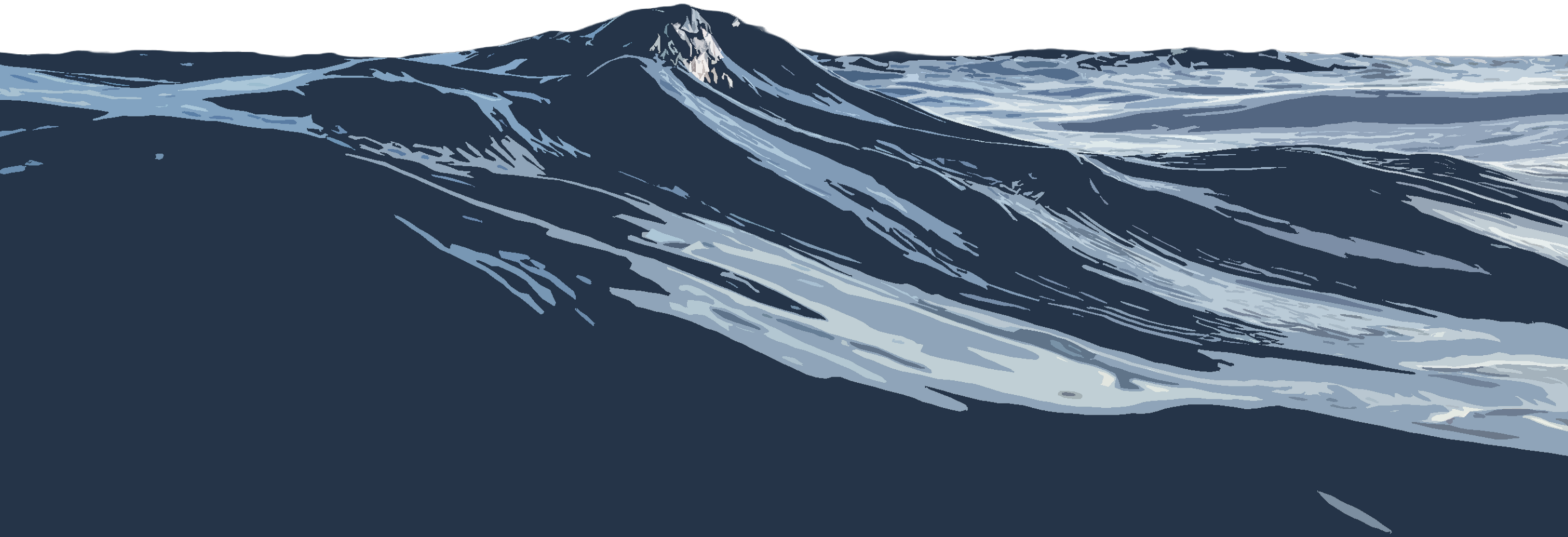


Conclusions

- **Bound harmonics** emerge with **increasing wind & wave steepness**, forming distinct nonlinear **dispersion branches**.
- Wave kinematics remain **consistent** with S&J throughout the **transition** of the underwater velocity profile to **turbulence**.
- Wave spectra exhibit significant **broadening** at fixed wavenumber, potentially encoding signatures of **turbulence fluctuations**.
- The subsurface **current profile** can be **obtained** by **inverting** the Doppler-shifted **dispersion relation**.
- The growth rate of higher harmonics can be written as $\beta_N(k) = N\beta_1(k/N)$
- Ongoing quantitative analysis of growth rate and broadening for various conditions

Introduction

Waves **exist** at the **interface** between air and water

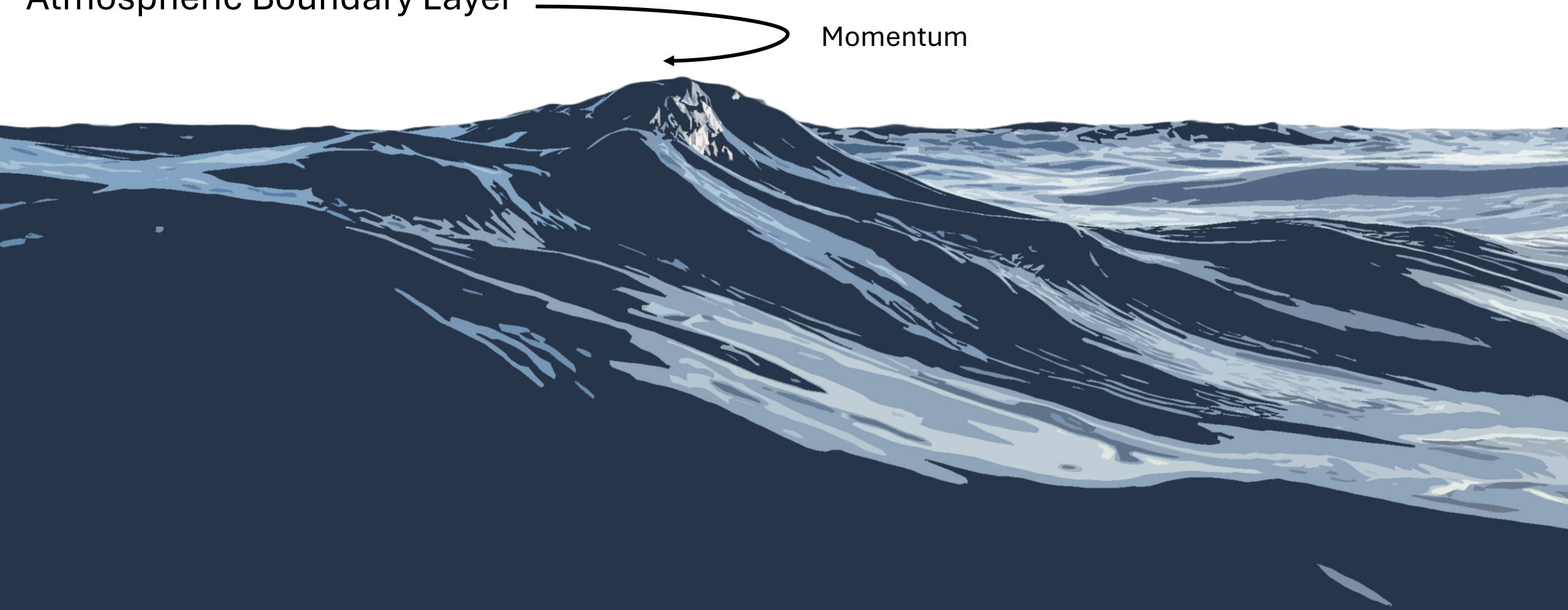


Introduction

Waves **exist** at the **interface** between air and water

Atmospheric Boundary Layer

Momentum

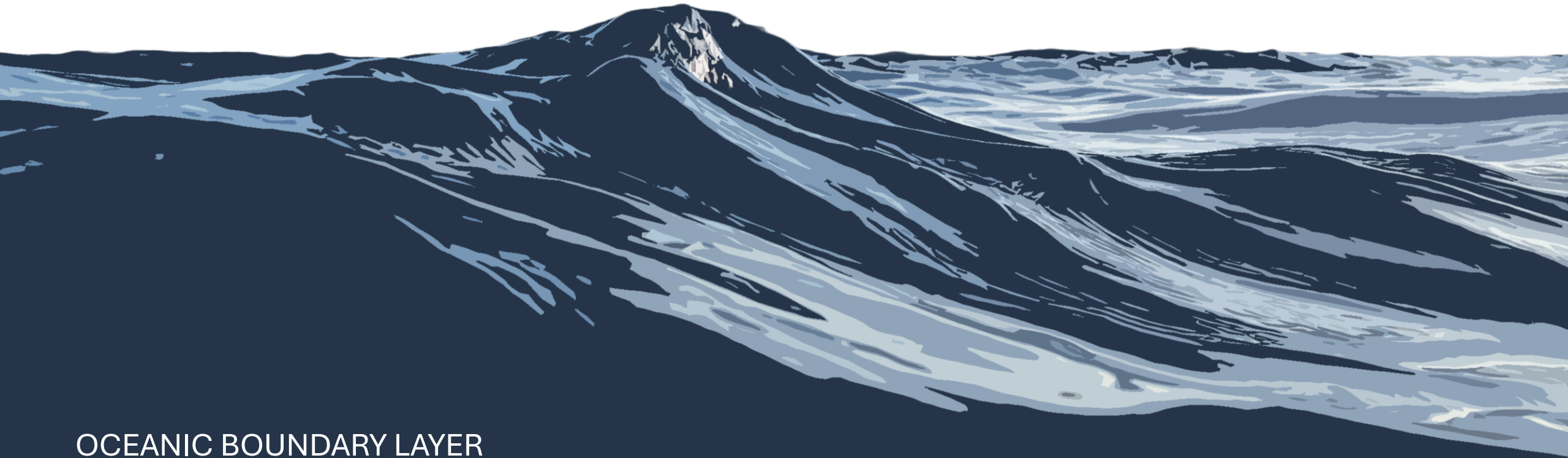


Introduction

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Atmospheric Boundary Layer

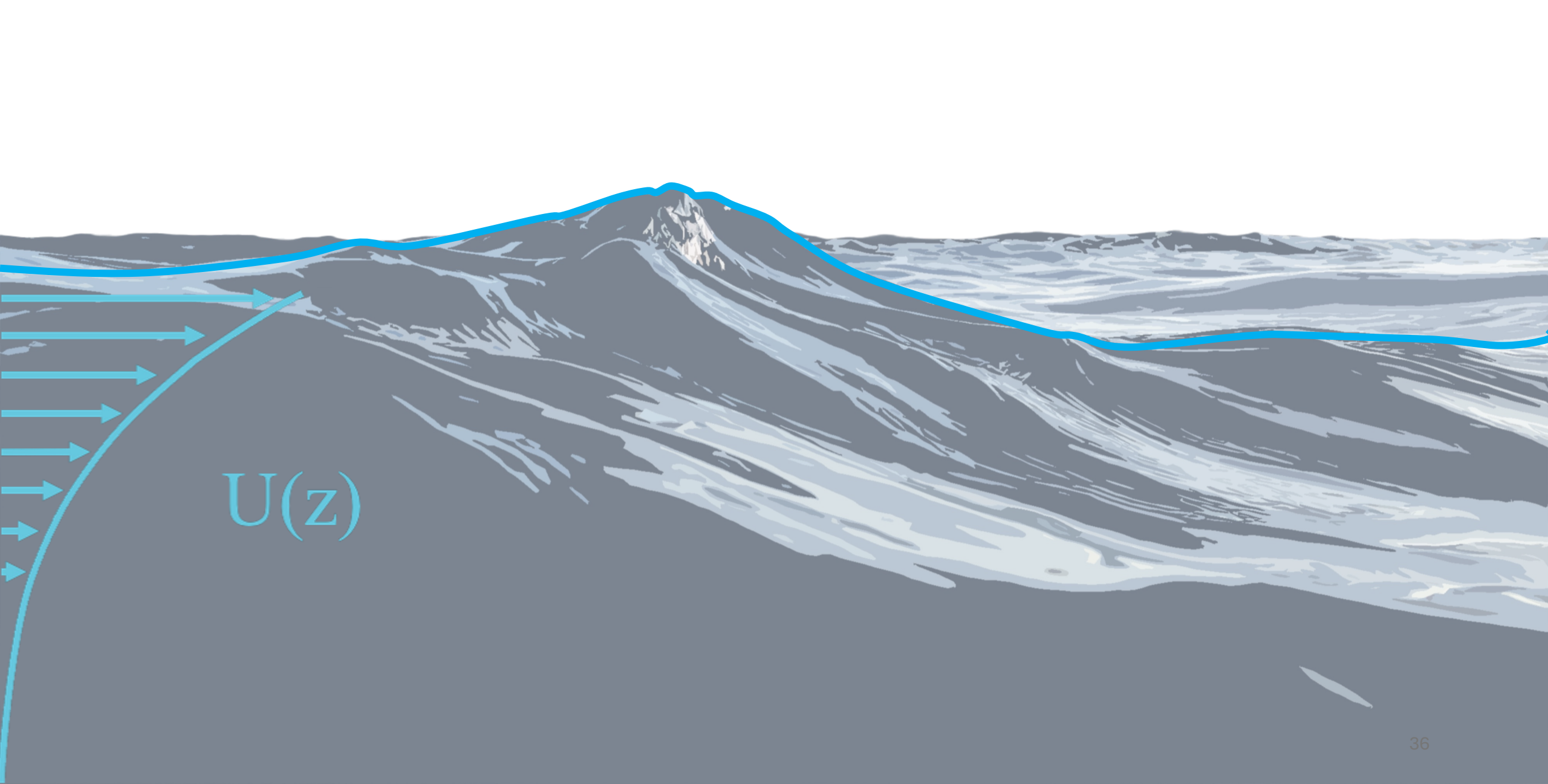
Momentum



OCEANIC BOUNDARY LAYER

Dissipate energy

Introduction

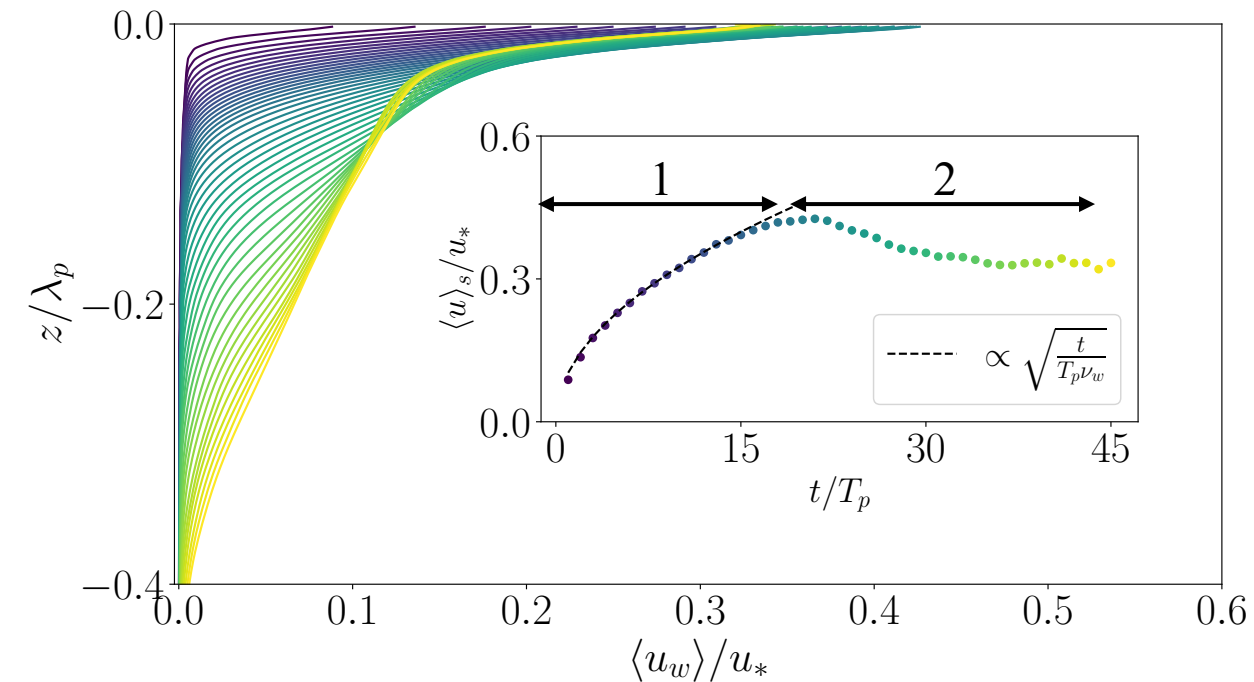
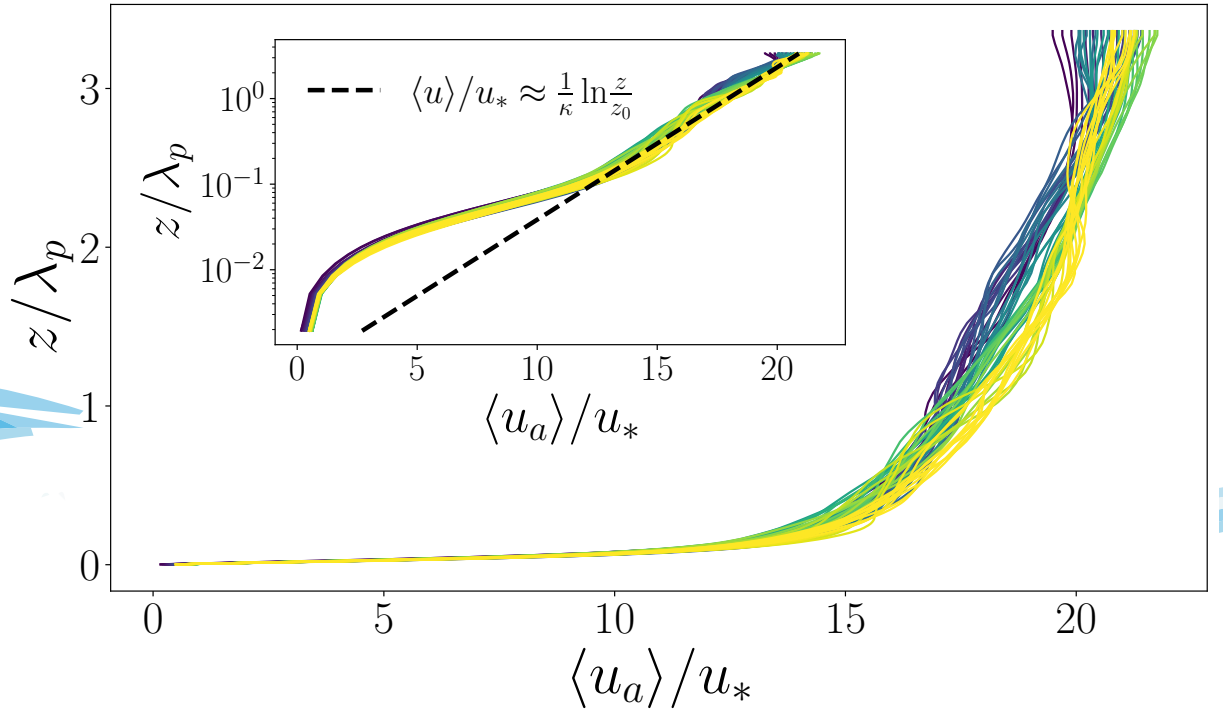
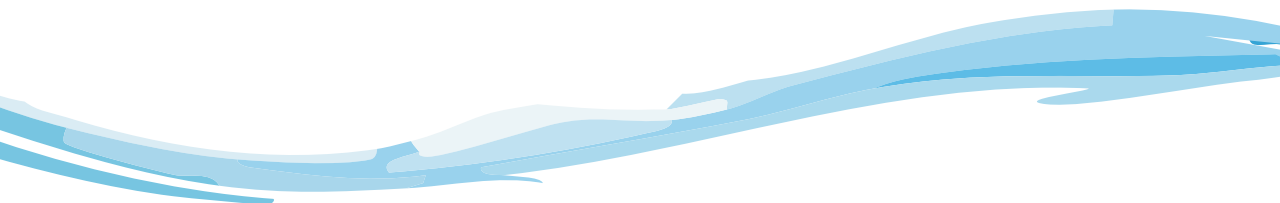


Mean Flow: Mean velocity profile in the air and the water

$$k_p H_s = 0.16$$

$$u_* / c = 0.5$$

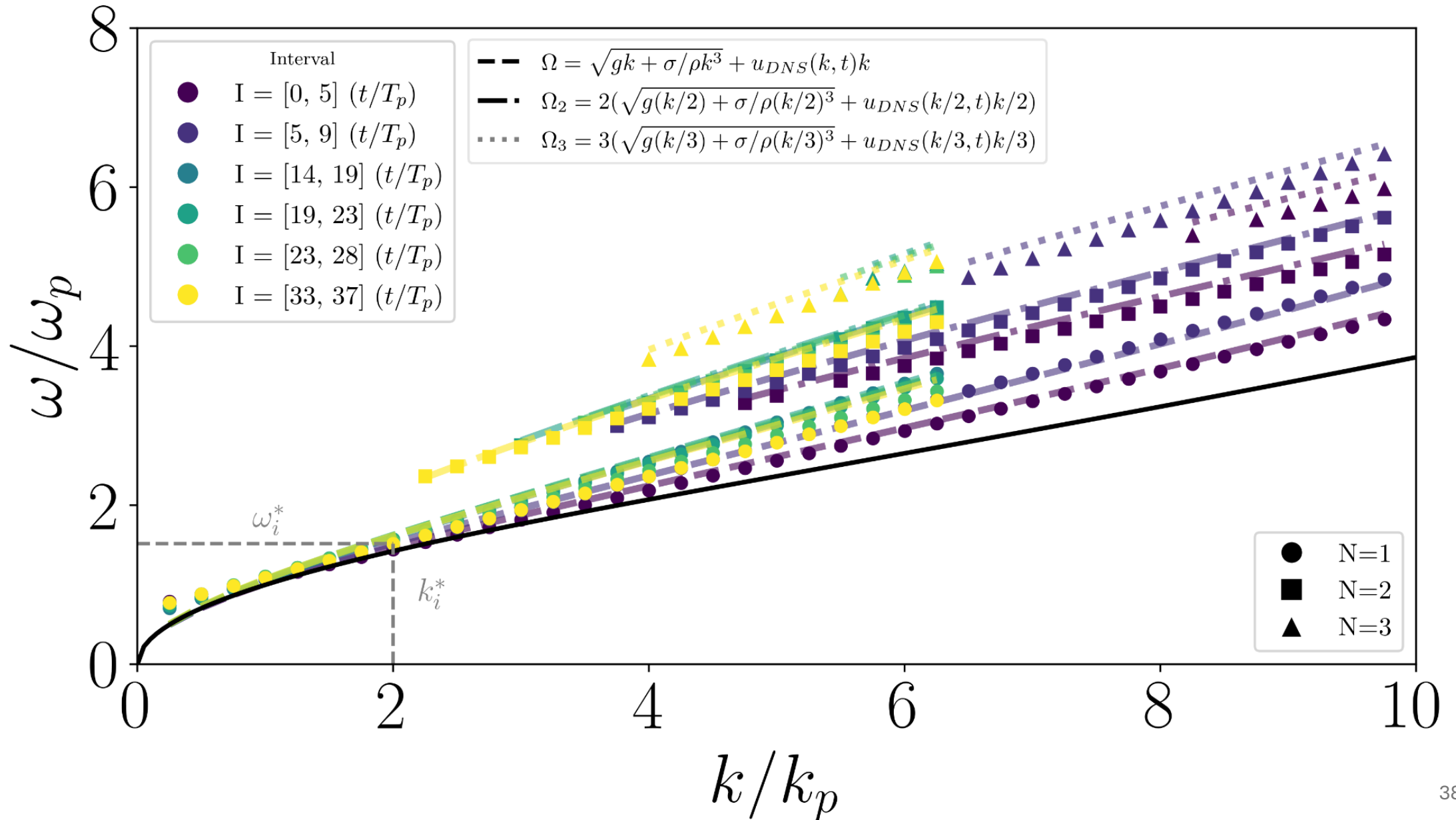
Air-side is **statistically stationary**



Developing viscous layer that transitions to a turbulent boundary layer

1. **Viscous** momentum diffusion
2. **Fully developed turbulence**

Phase speed of the higher bound modes



Branches Time Evolution: Higher harmonics and non-linear dispersion relation

Primary mode

$$(k_*, \omega_*)$$

Non-linear interaction with itself

Higher harmonics

$$\begin{array}{l} (2k_*, 2\omega_*) \\ (3k_*, 3\omega_*) \end{array} \left| \begin{array}{l} k_N = Nk_* \\ \Omega_N = N\omega_* \end{array} \right.$$

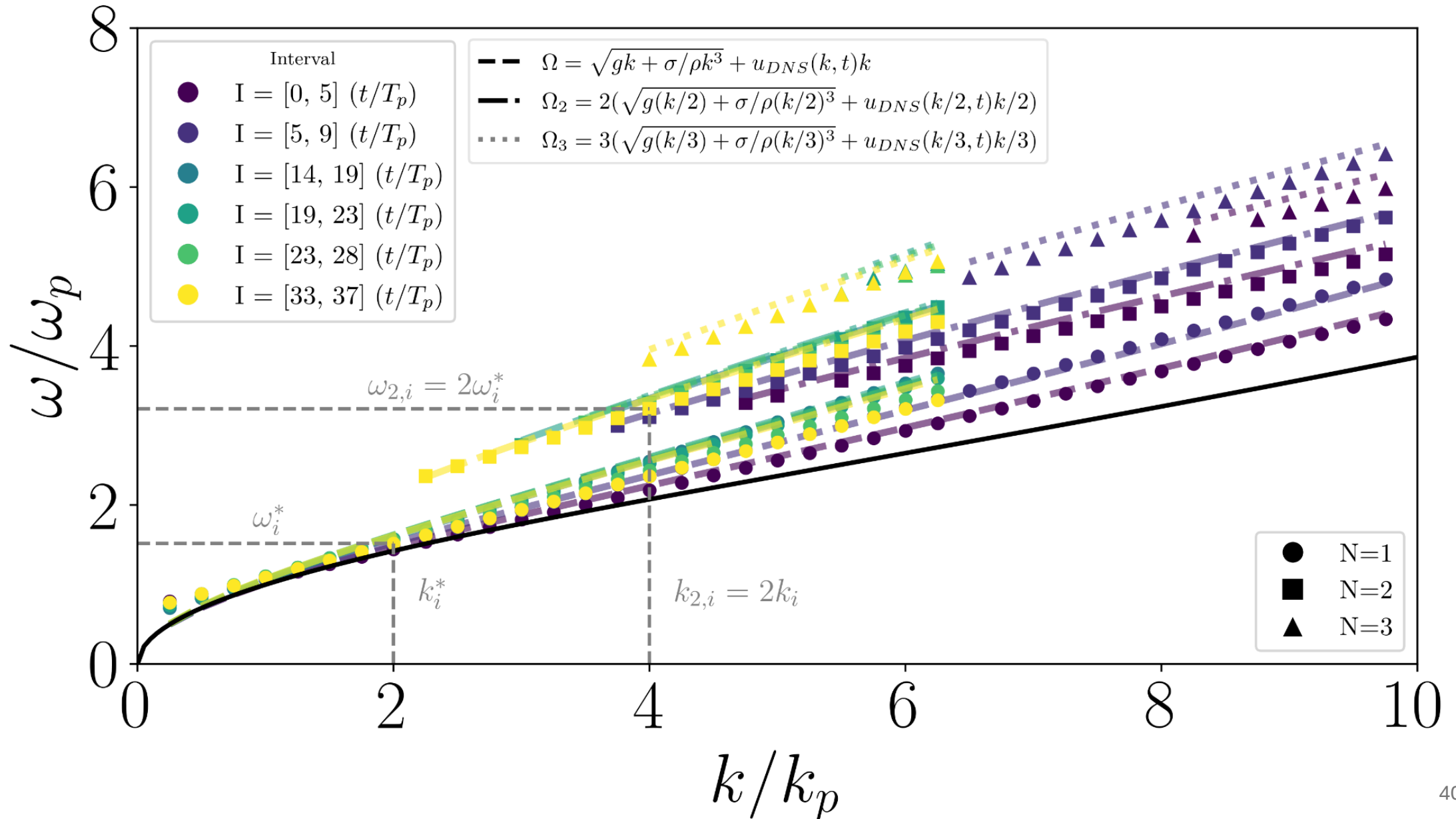
$$\Omega_N(k_N) = N \sqrt{\frac{gk_N}{N} + \frac{\sigma}{\rho} \left(\frac{k_N}{N}\right)^3}$$

Non-linear dispersion relation

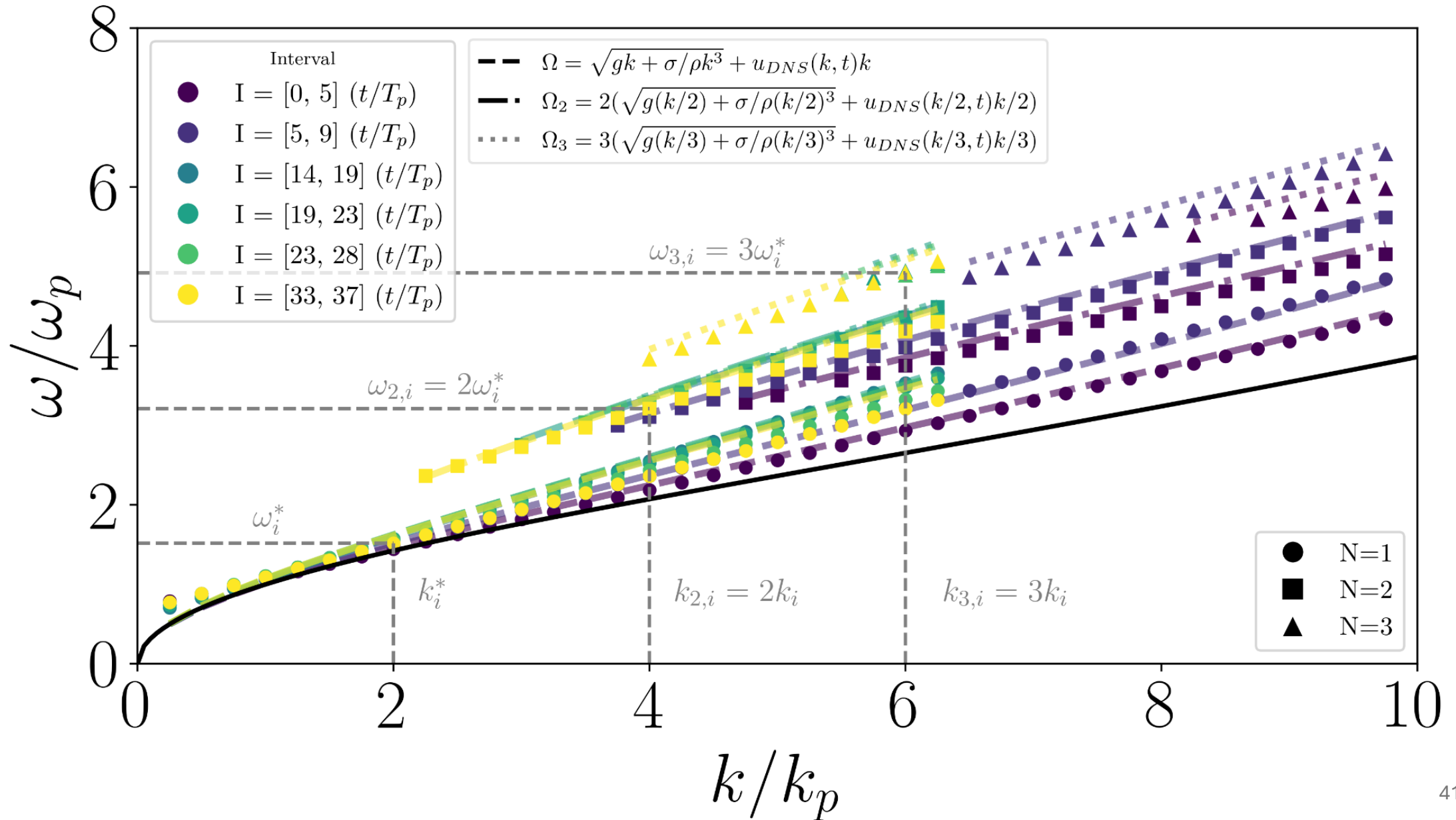
$$\Omega_N(k_N) = N \left[\sqrt{\frac{gk_N}{N} + \frac{\sigma}{\rho} \left(\frac{k_N}{N}\right)^3} + \frac{u_{eff} \left(\frac{k_N}{N}, t\right) k_N}{N} \right]$$

with the Doppler shift from Stewart & Joy (1974) $u_{eff}(k) = 2k \int_0^\infty u_L(z) e^{-2kz} dz$

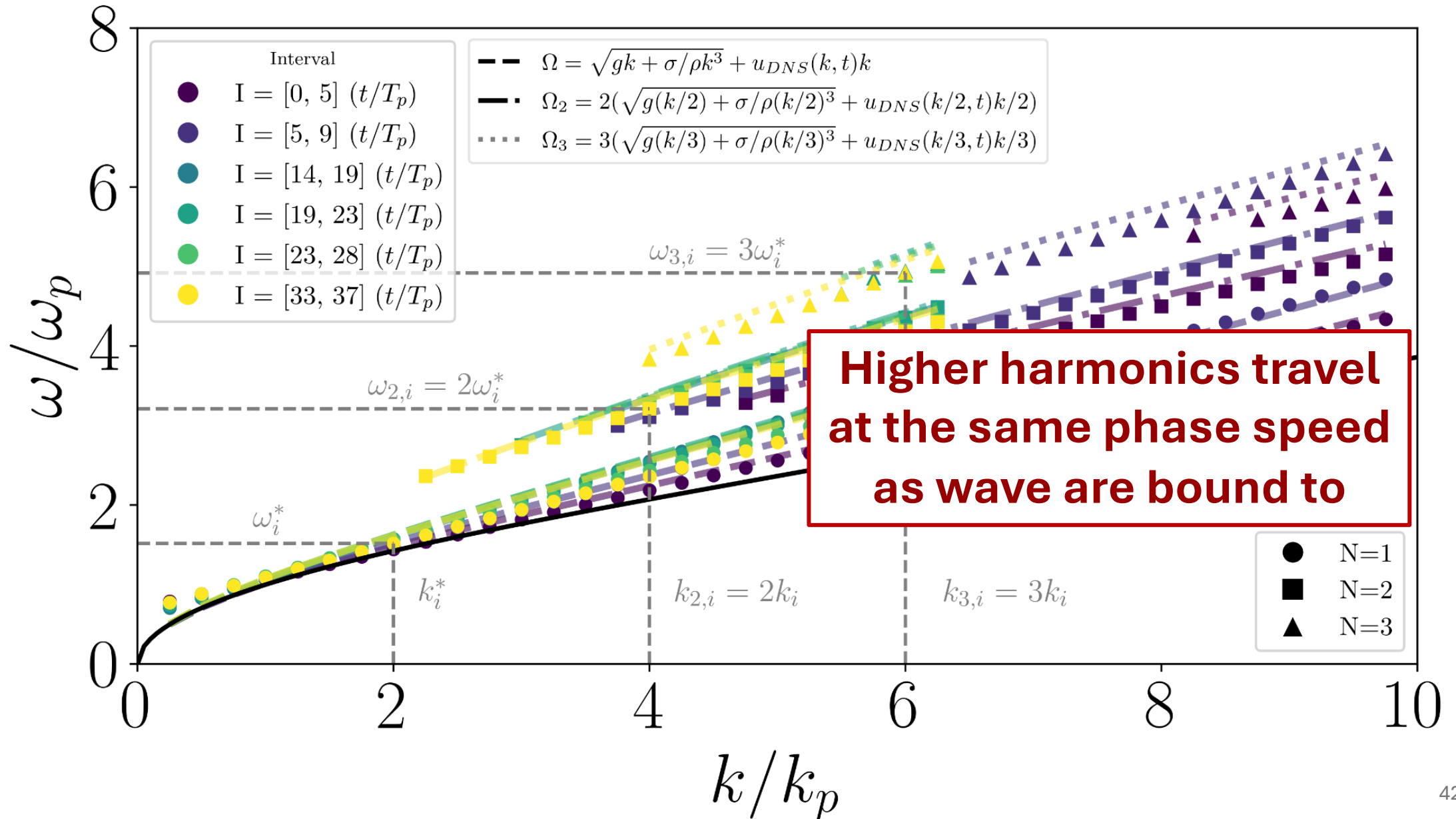
Phase speed of the higher bound modes



Phase speed of the higher bound modes



Phase speed of the higher bound modes



Stability and convergence

RMS surface elevation

$$\sigma_\eta = \sqrt{\langle \eta^2 \rangle}$$

Viscous length

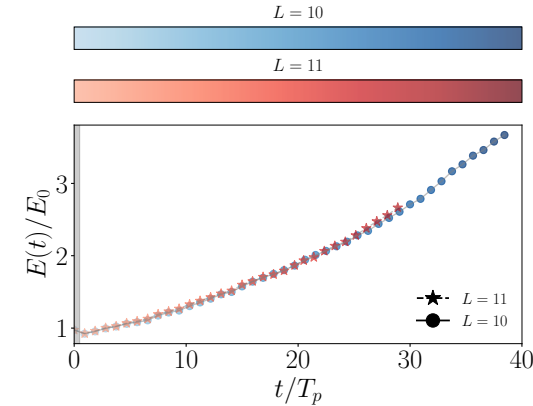
$$\delta_\nu = \lambda_p / \sqrt{Re}$$

Kolmogorov length scale

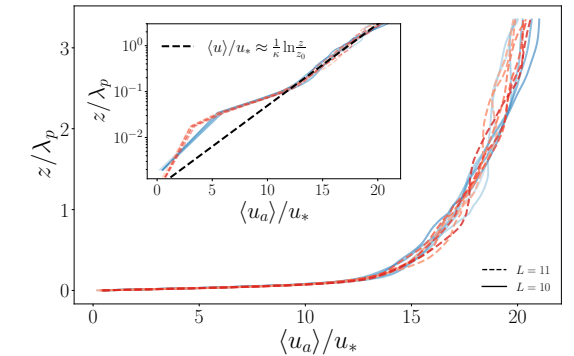
$$\eta_k = \left(\frac{\nu_w^3}{\epsilon_{max}} \right)^{1/4} \quad \epsilon = 2\nu_w S_{ij} S_{ij} \quad S_{ij} = \frac{1}{2} \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$$

$k_p H_s$	u_*/c	Bo	L	σ_η/Δ	δ_ν/Δ	η_K/Δ
0.04-0.16	0.25-0.50	25, 200, 1000	10	1-2.3	1.6	2.2-2.6
0.16	0.50	200	11	1.5-4.2	3.2	2.2-2.7

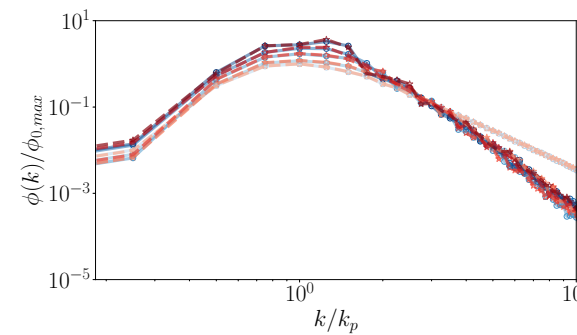
a) Potential wave energy



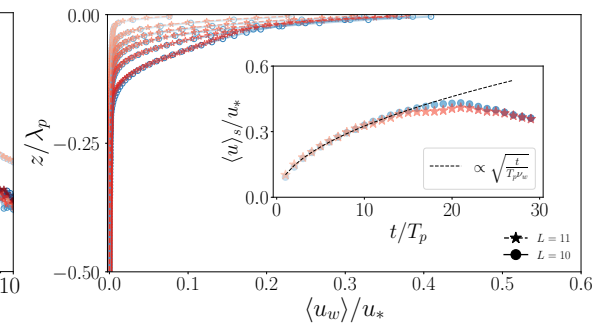
c) Air-side mean velocity profiles



b) Wave energy spectrum

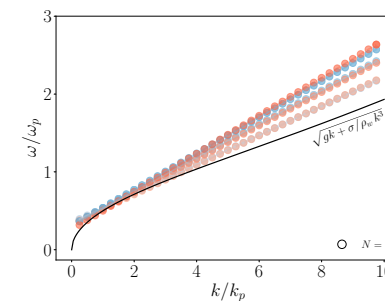


d) Water-side mean velocity profiles

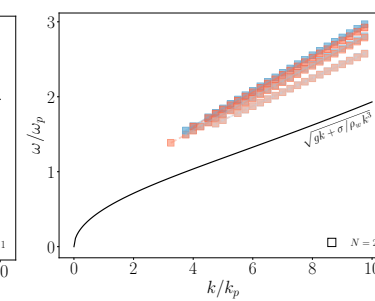


Doppler shifted linear dispersion relation

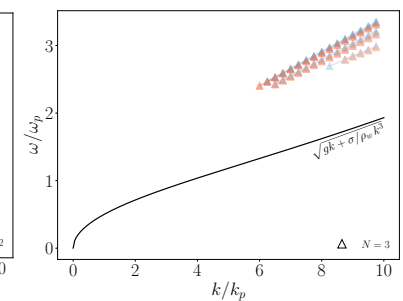
e) Main harmonic



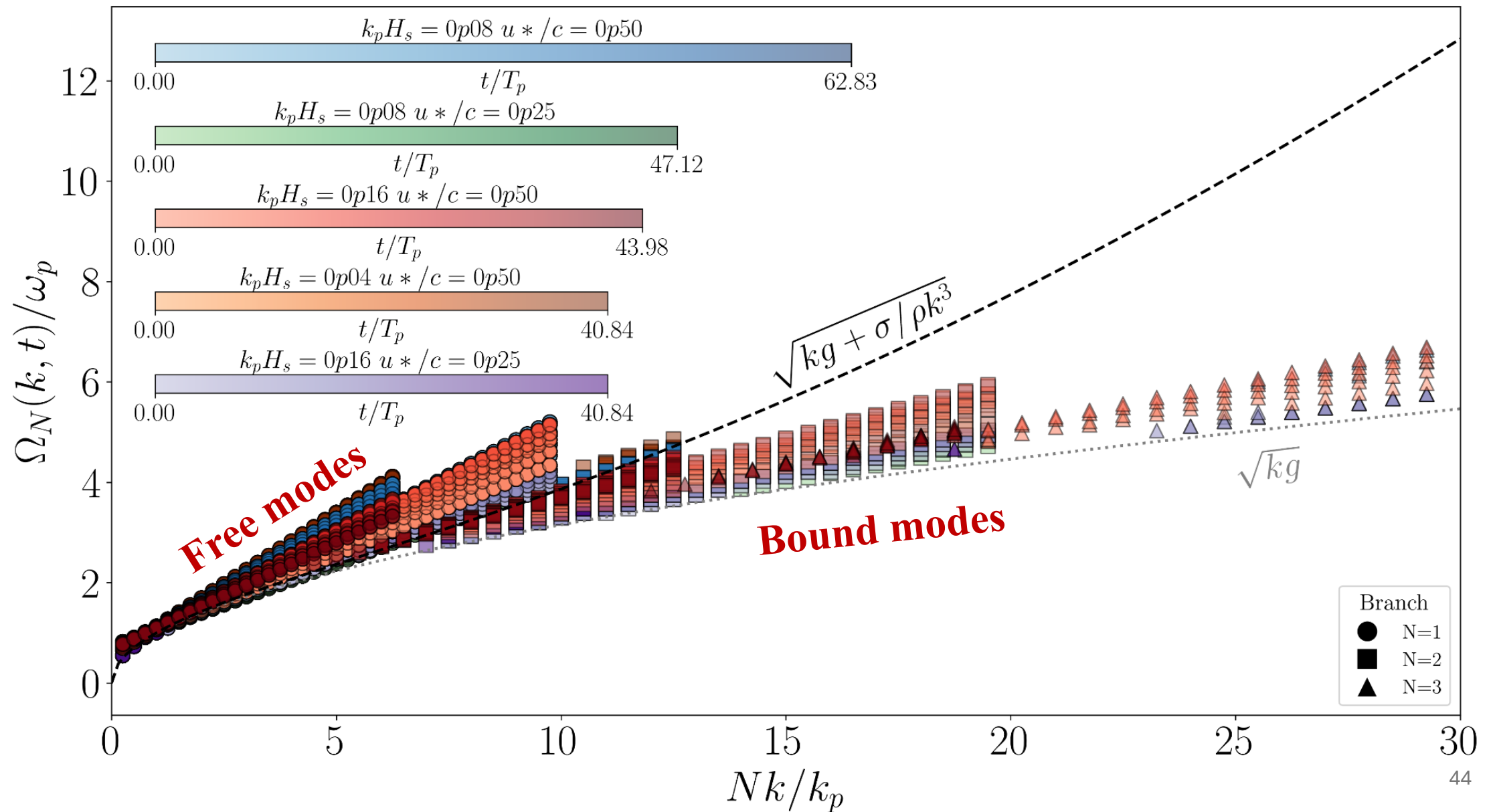
f) Second harmonic



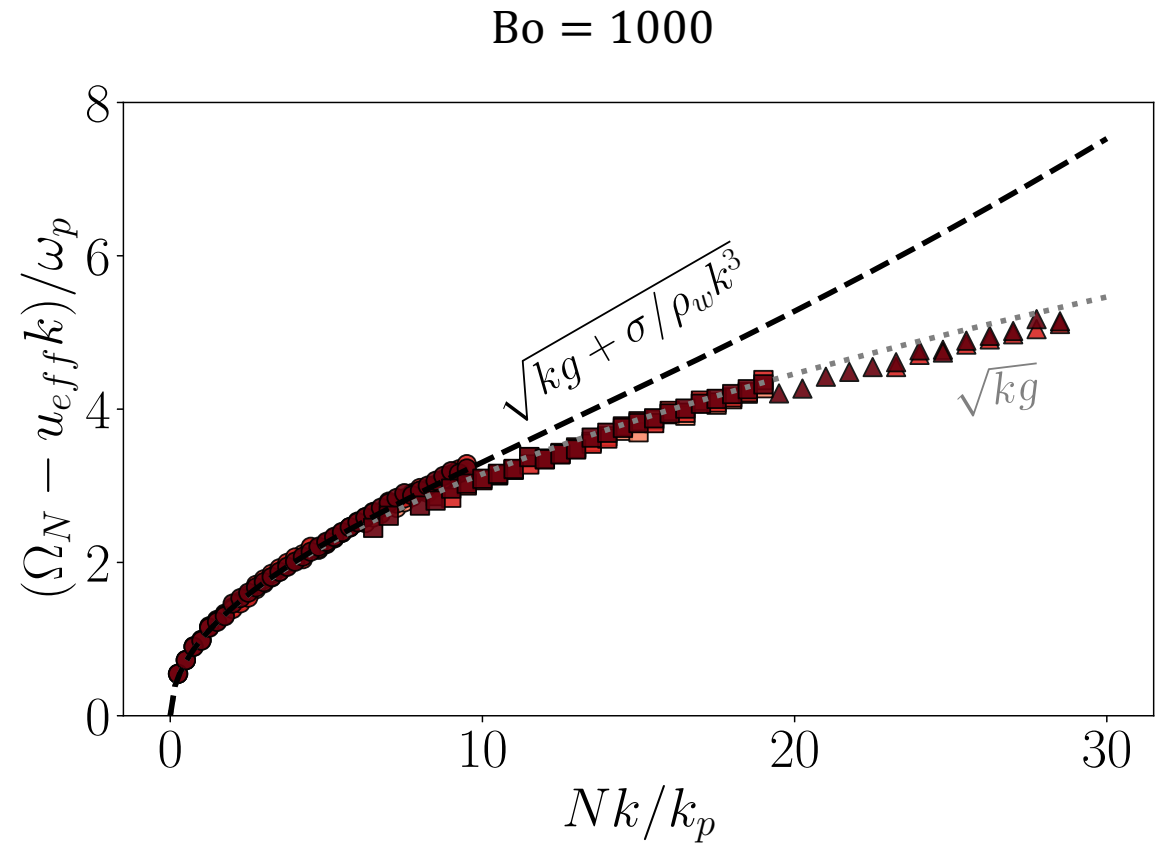
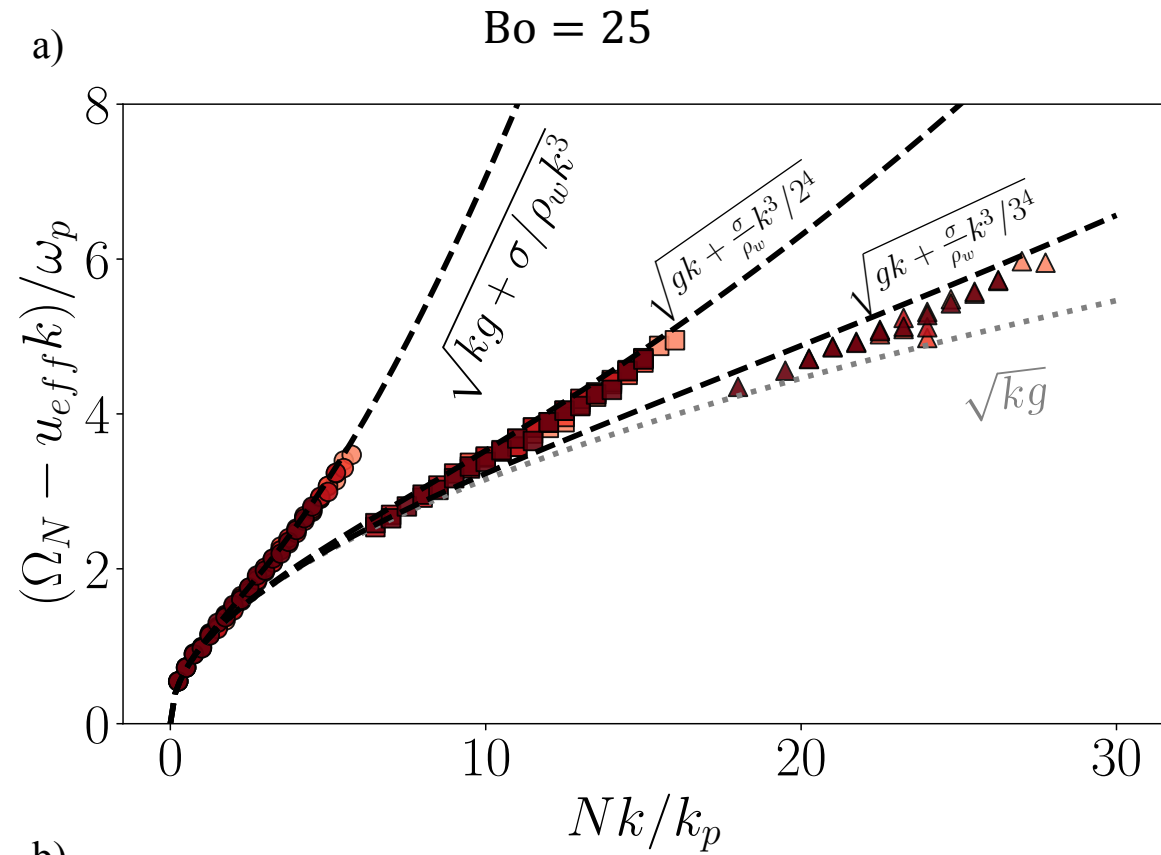
g) Third harmonic



Does the non-linear dispersion relation hold across different wind and wave initial conditions?

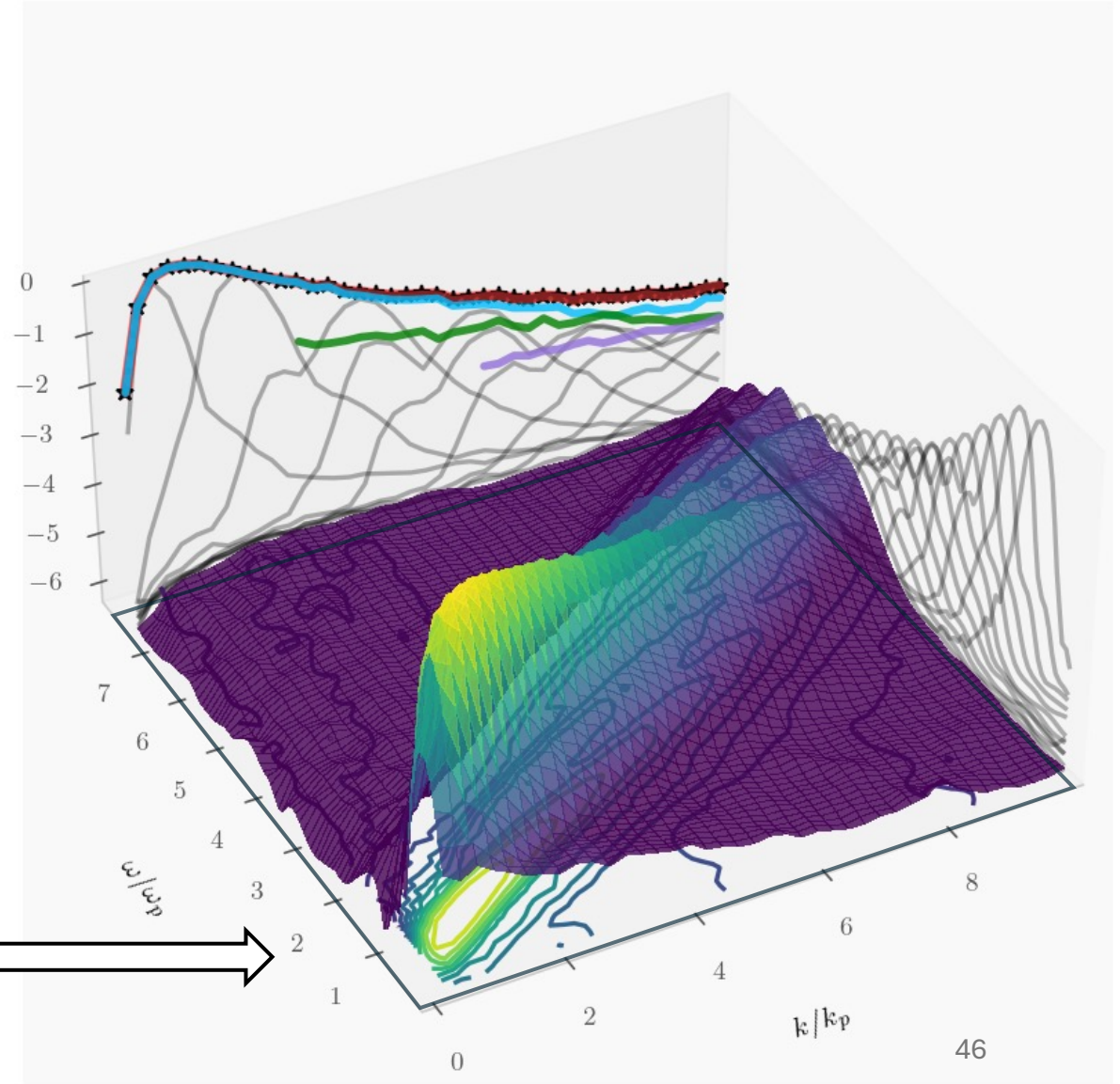
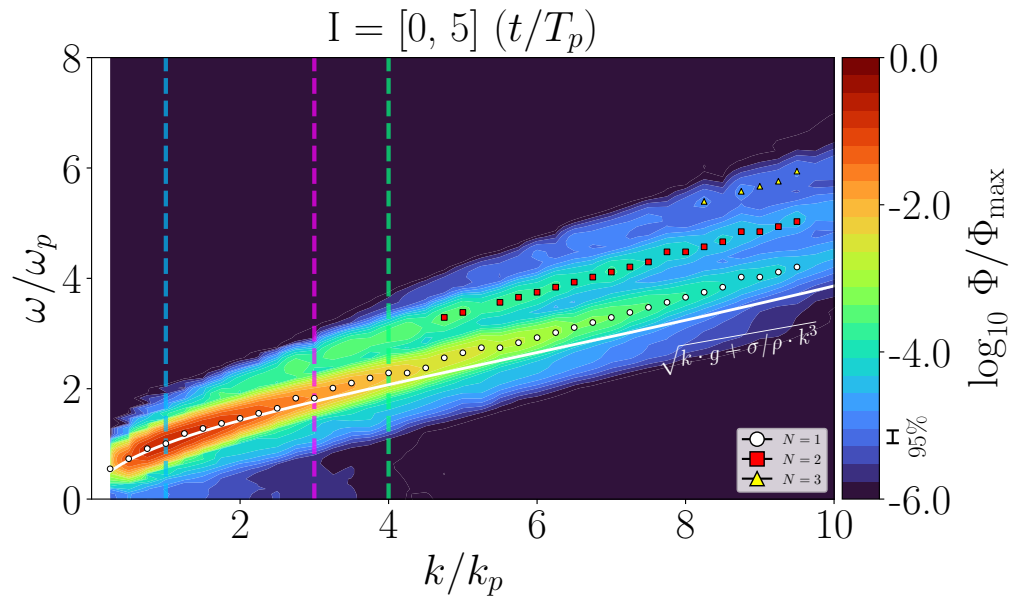
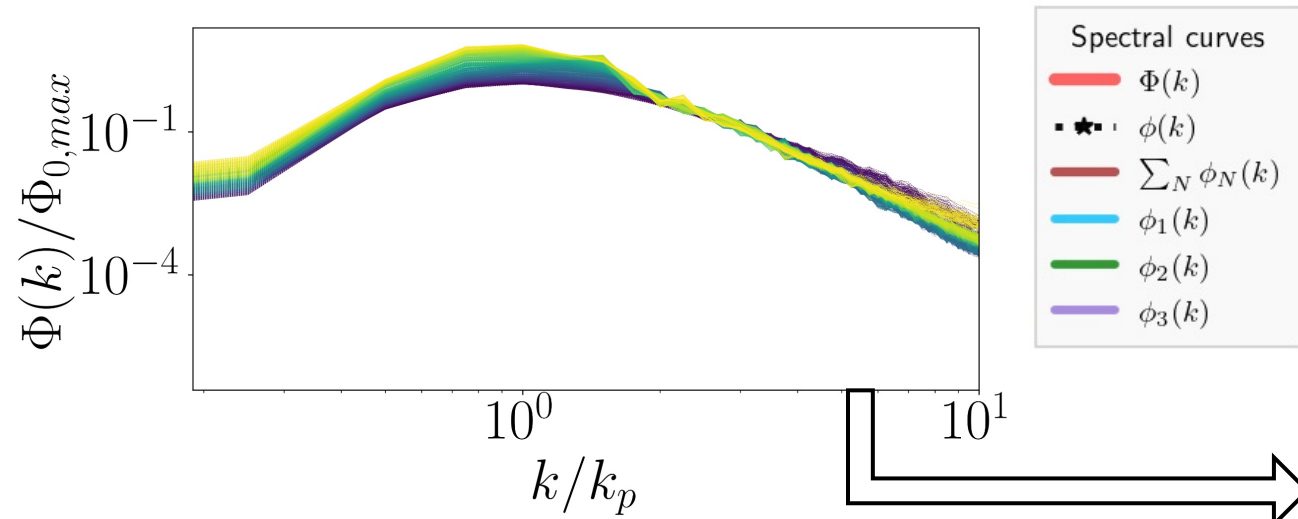


Different gravity-capillary regimes

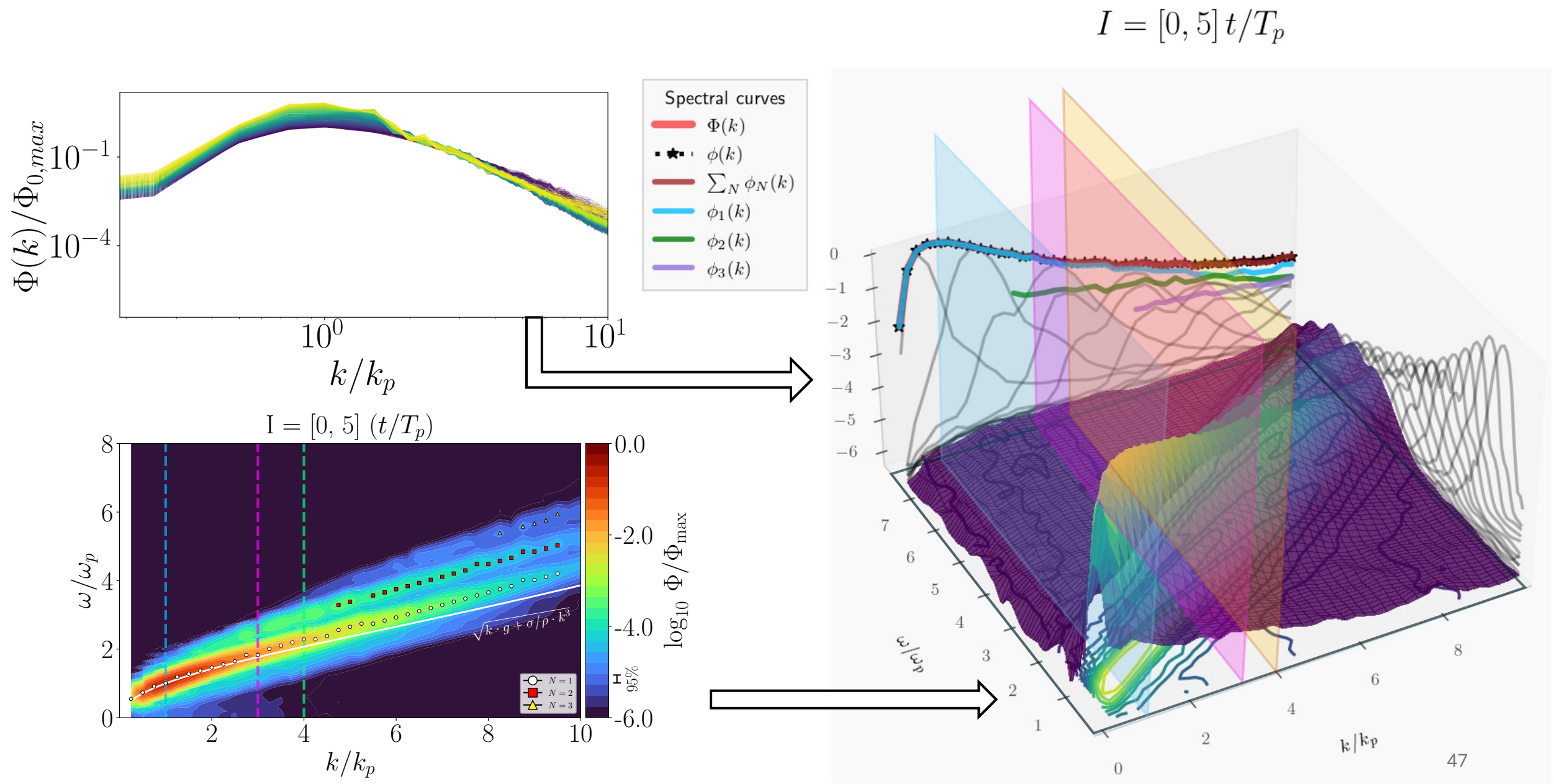


$(k - \omega)$ spectral analysis: growth of modes $\beta_\Phi(k)$ and $\beta_N(k)$

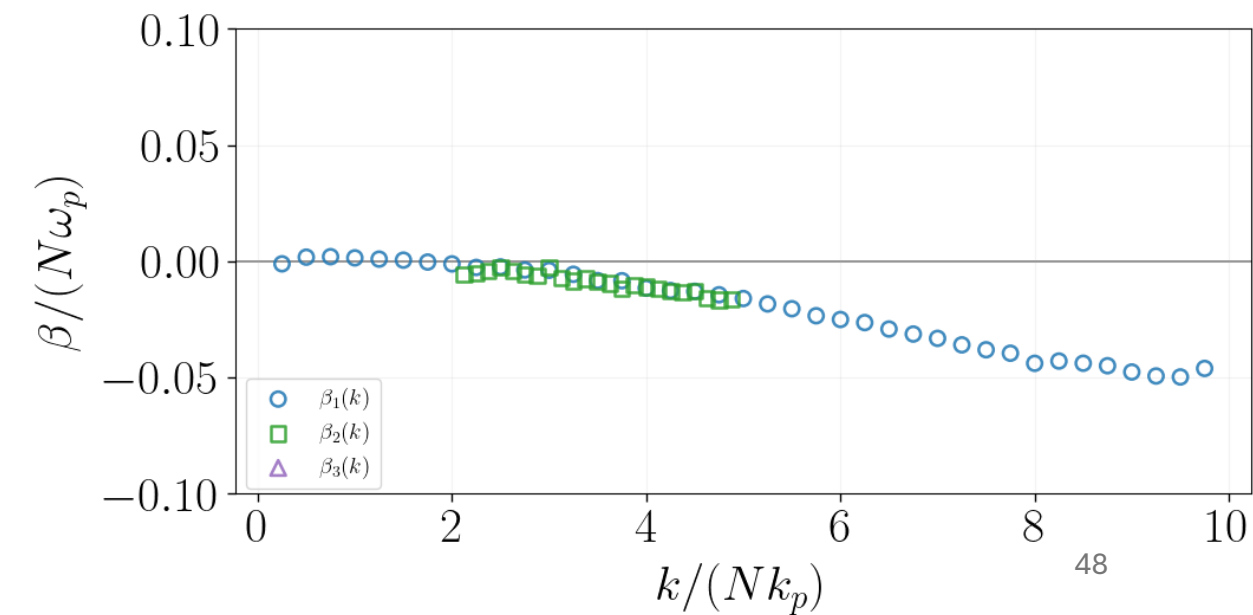
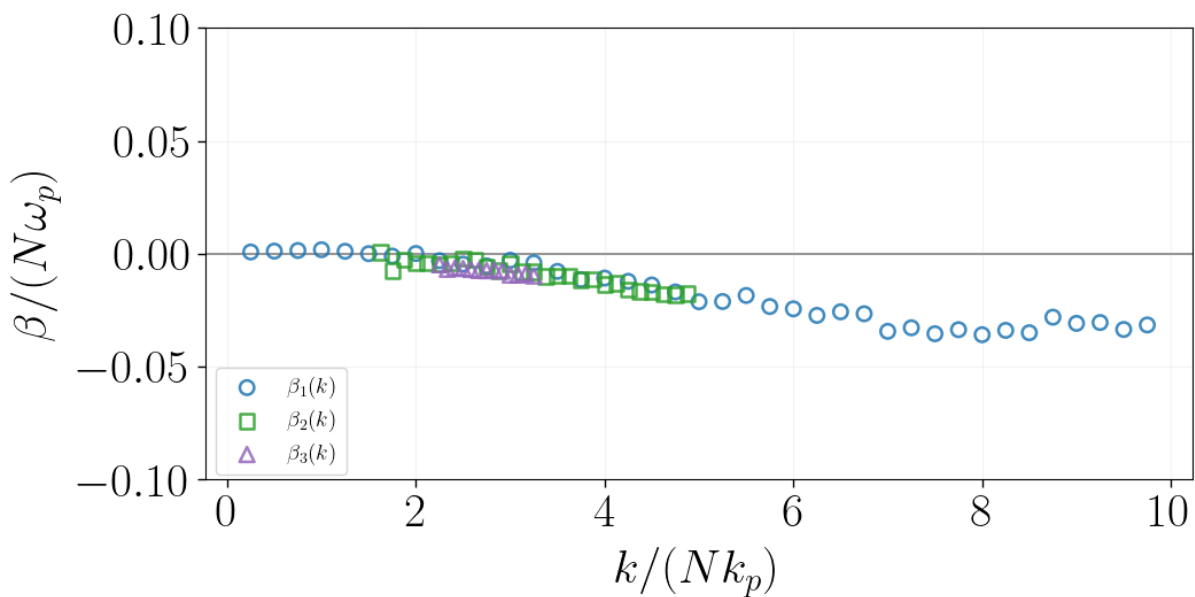
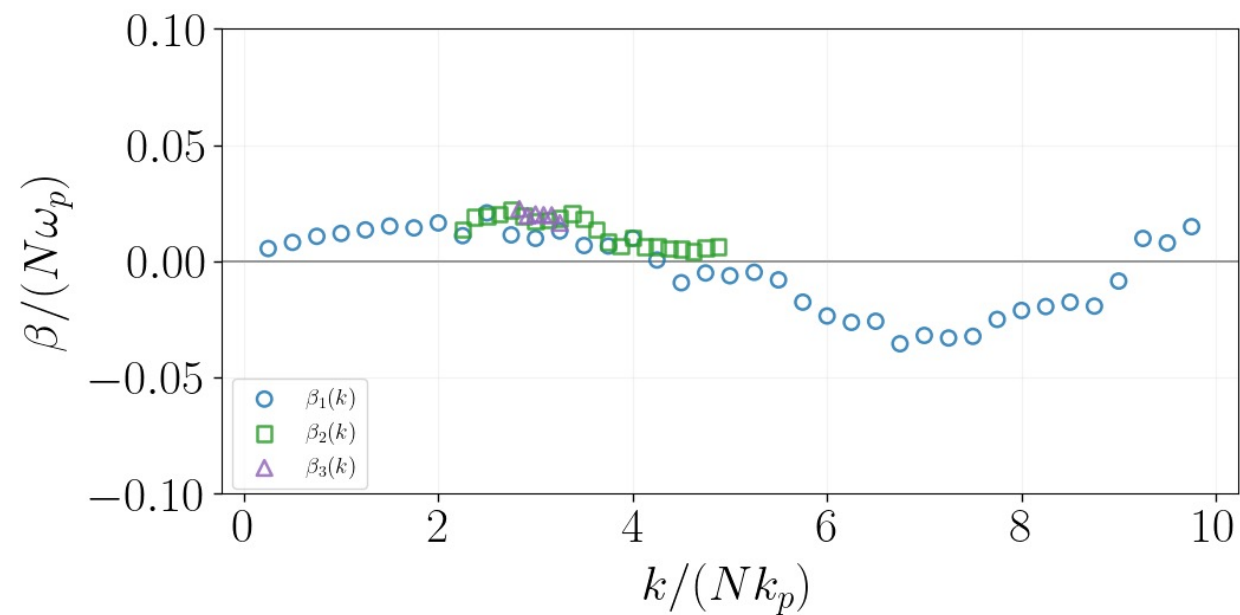
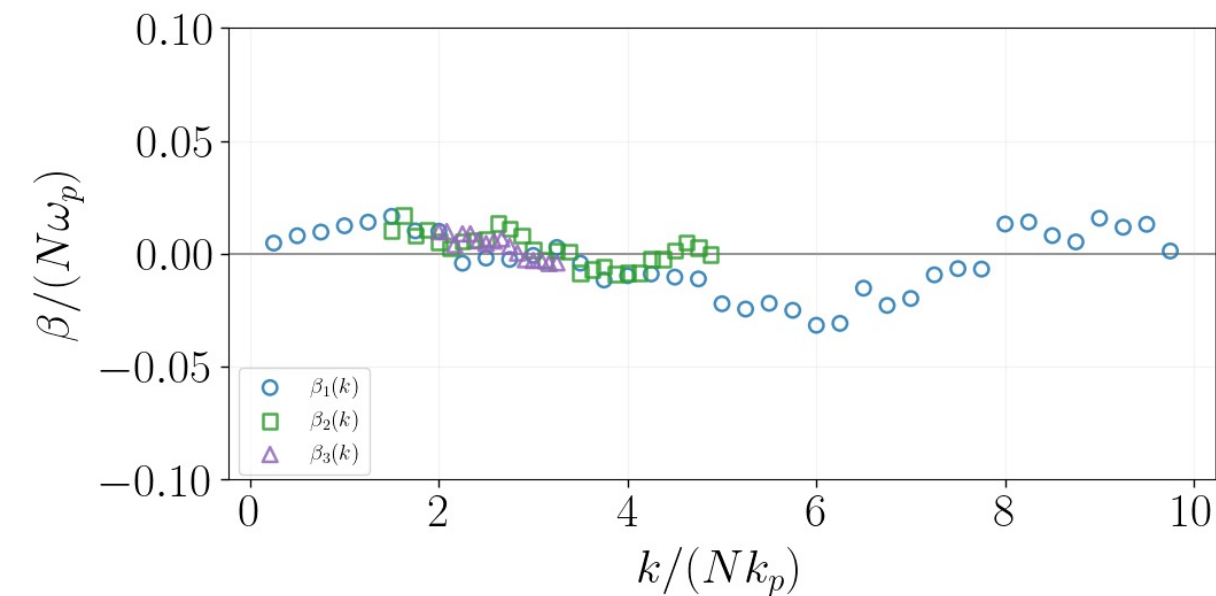
$I = [0, 5] t/T_p$



$(k - \omega)$ spectral analysis: growth of modes $\beta_\Phi(k)$ and $\beta_N(k, N)$



Energy Extraction: Growth rate per branch accounting for Doppler shift $\beta_N(k, N)$



Can we reconstruct $\beta(k)$ from $\beta_N(k, N)$ of each branch?

k space

$\phi(k)$

Growth rate $\beta(k)$

$$\beta(k) = \frac{d}{dt} \ln \phi(k, t) = \frac{1}{\phi} \frac{\partial \phi}{\partial t}$$

$$\Phi(k) = \sum_N \Phi_N(k) = \phi(k)$$

$$\beta(k) = \frac{1}{\sum_N \Phi_N} \sum_N \frac{\partial \Phi_N}{\partial t}$$

$(k - \omega)$ space

$$\Phi(k, t) = \int \Phi(k, \omega', t) d\omega' \approx \underbrace{\sum_N \int_{B_N} \Phi d\omega'}_{\text{branches}} = \sum_N \Phi_N(k, t) \quad \text{where} \quad \Phi_N(k, t) = \int_{\omega \in B_N} \Phi(k, \omega', t) d\omega'$$

Growth rate $\beta_N(k)$

$$\beta_N(k) = \frac{d}{dt} \ln \Phi_N(k, t) = \frac{1}{\Phi_N} \frac{\partial \Phi_N}{\partial t} \longrightarrow \frac{\partial \Phi_N}{\partial t} = \beta_N \Phi_N$$

$$\beta(k) = \frac{\sum_N \beta_N \Phi_N}{\sum_N \Phi_N}$$

Weighted average

Can we reconstruct $\beta(k)$ from $\beta_N(k, N)$ of each branch?

k space

$\phi(k)$

Growth rate $\beta(k)$

$$\beta(k) = \frac{d}{dt} \ln \phi(k, t) = \frac{1}{\phi} \frac{\partial \phi}{\partial t}$$

$$\Phi(k) = \sum_N \Phi_N(k) = \phi(k)$$

$$\beta(k) = \frac{1}{\sum_N \Phi_N} \sum_N \frac{\partial \Phi_N}{\partial t}$$

$(k - \omega)$ space

$$\Phi(k, t) = \int \Phi(k, \omega', t) d\omega' \approx \underbrace{\sum_N \int_{B_N} \Phi d\omega'}_{\text{branches}} = \sum_N \Phi_N(k, t) \quad \text{where} \quad \Phi_N(k, t) = \int_{\omega \in B_N} \Phi(k, \omega', t) d\omega'$$

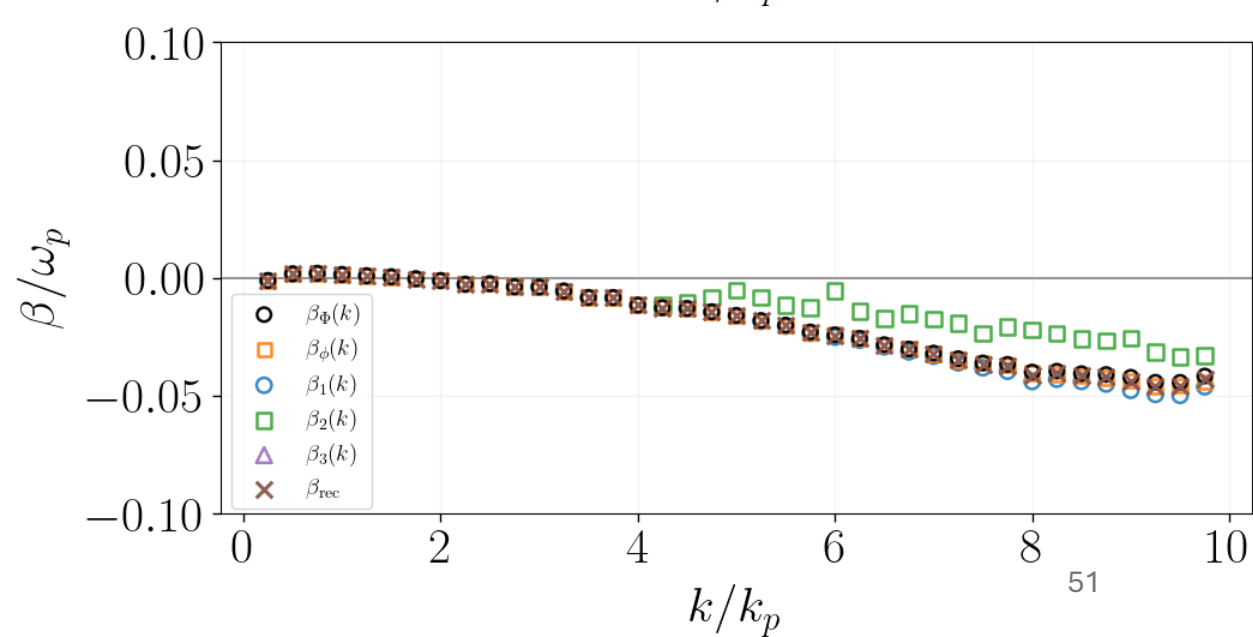
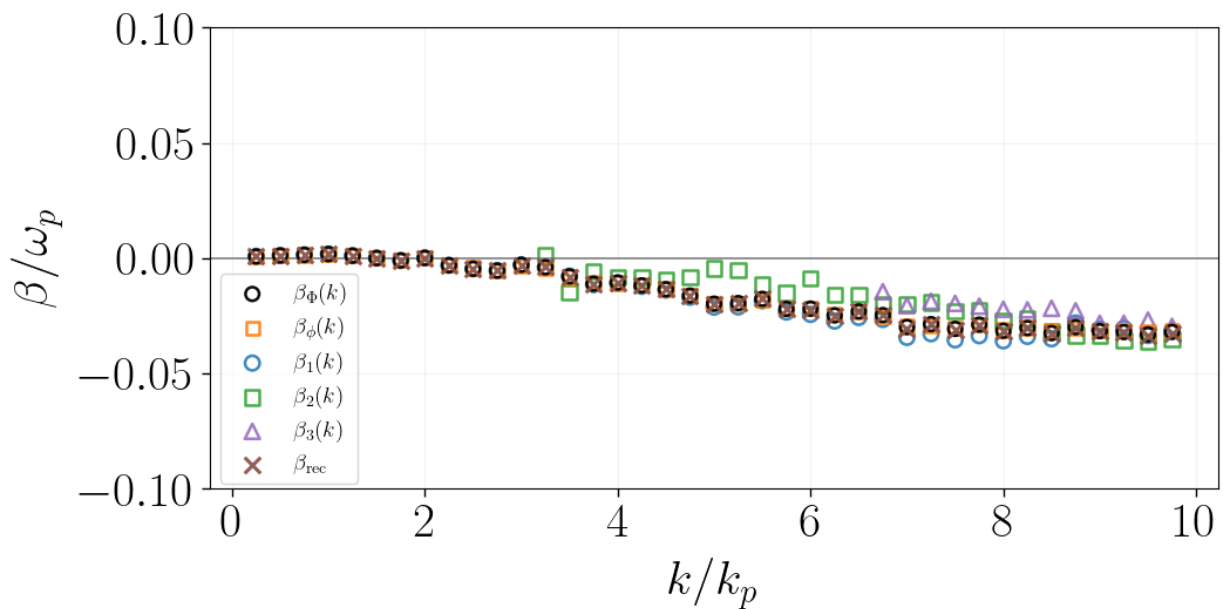
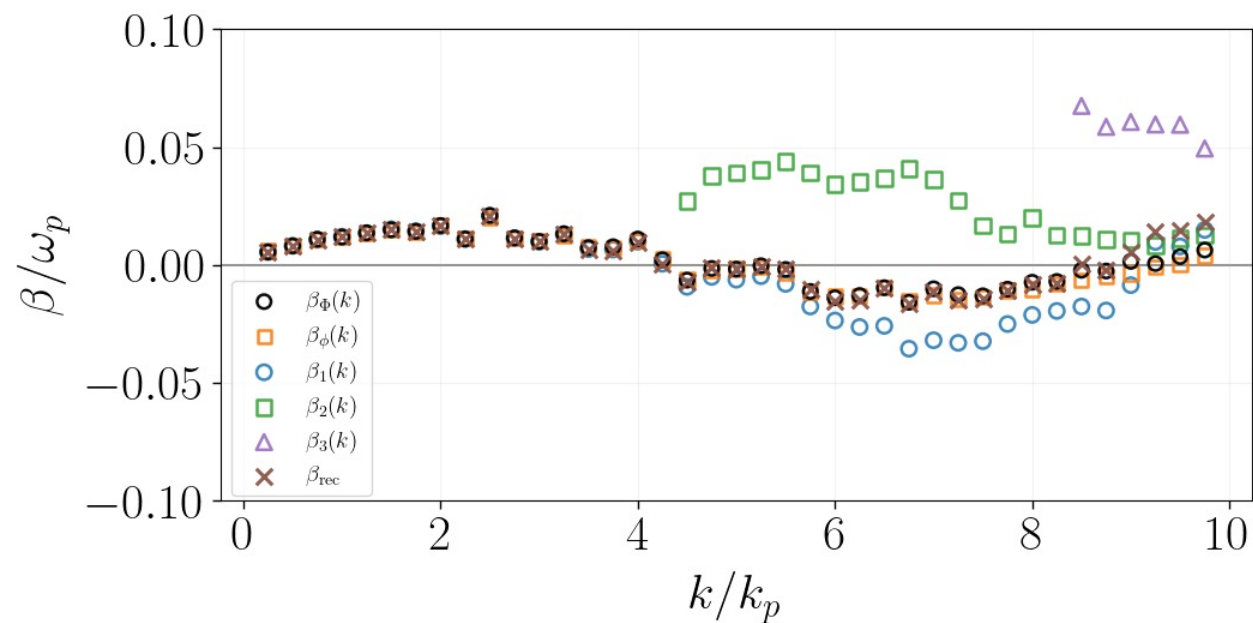
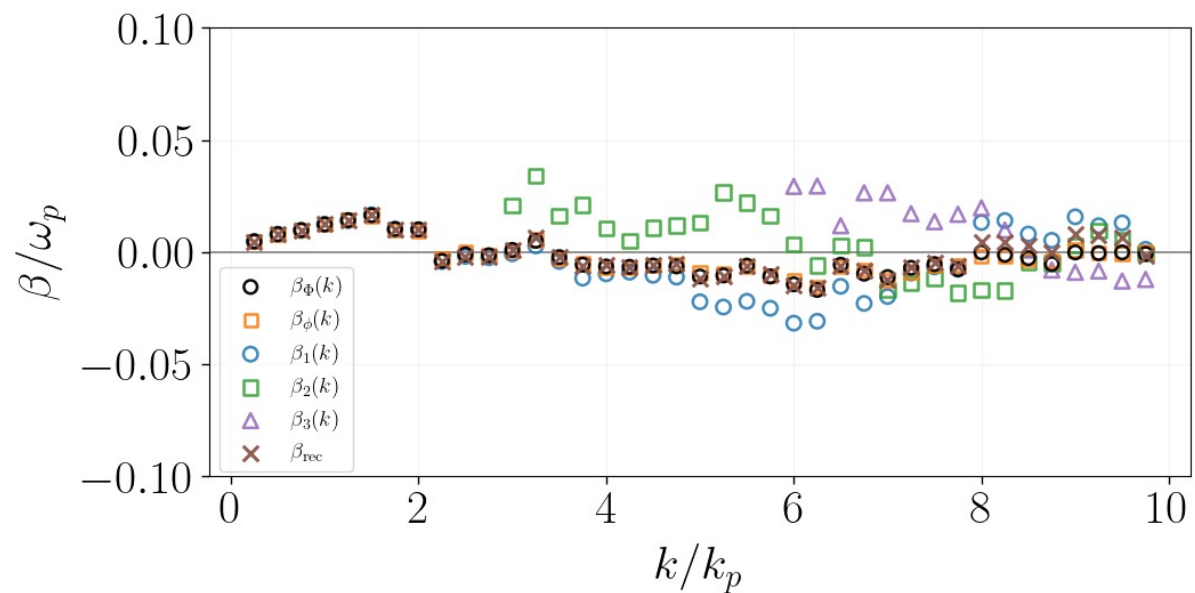
branches

Growth rate $\beta_N(k)$

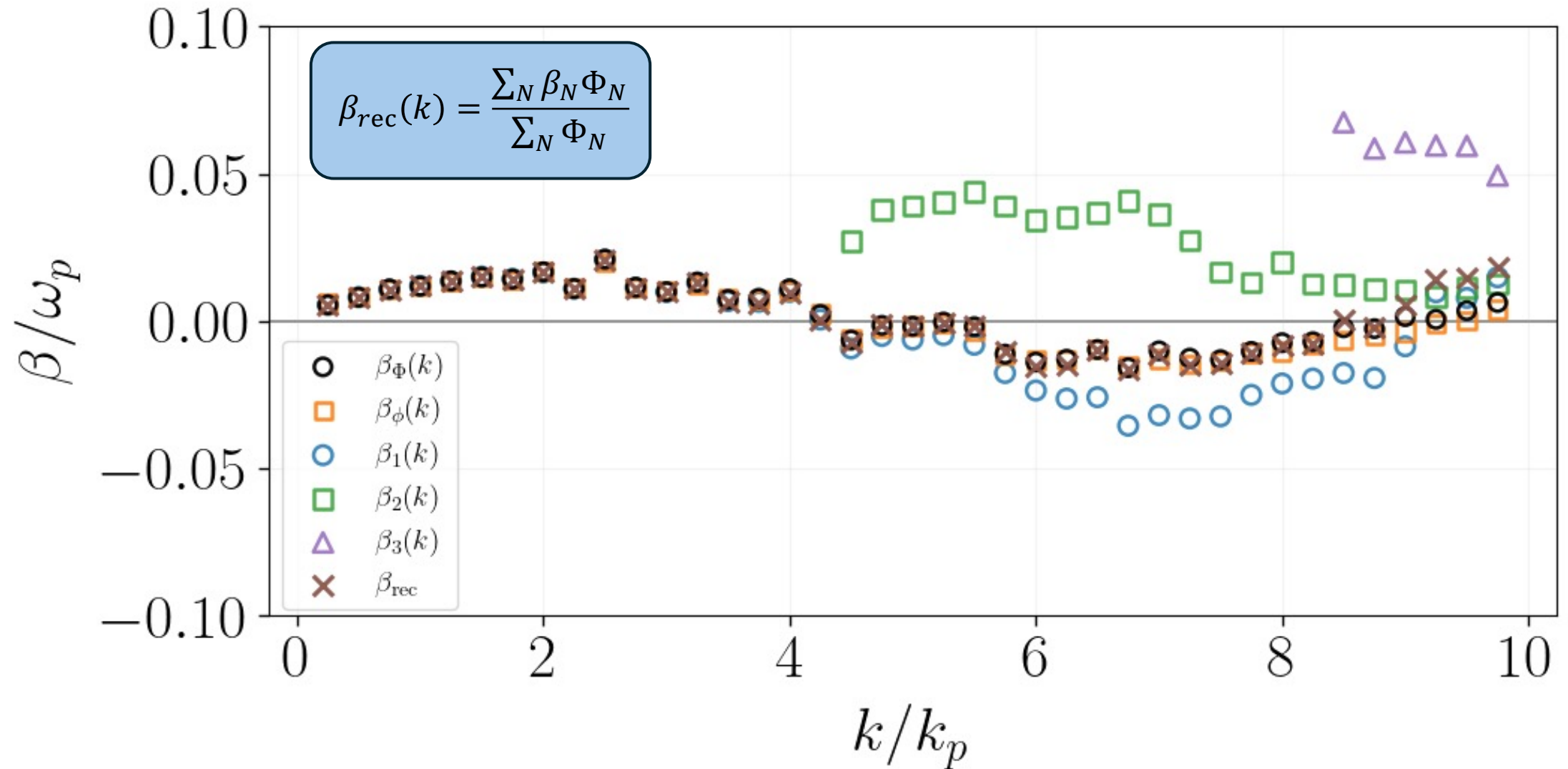
$$\beta_N(k) = \frac{d}{dt} \ln \Phi_N(k, t) = \frac{1}{\Phi_N} \frac{\partial \Phi_N}{\partial t} \longrightarrow \frac{\partial \Phi_N}{\partial t} = \beta_N \Phi_N$$

$$\beta_{\text{rec}}(k) = \frac{\sum_N \beta_N \Phi_N}{\sum_N \Phi_N}$$

Energy Extraction: Growth rate per branch accounting for Doppler shift $\beta_N(k, N)$



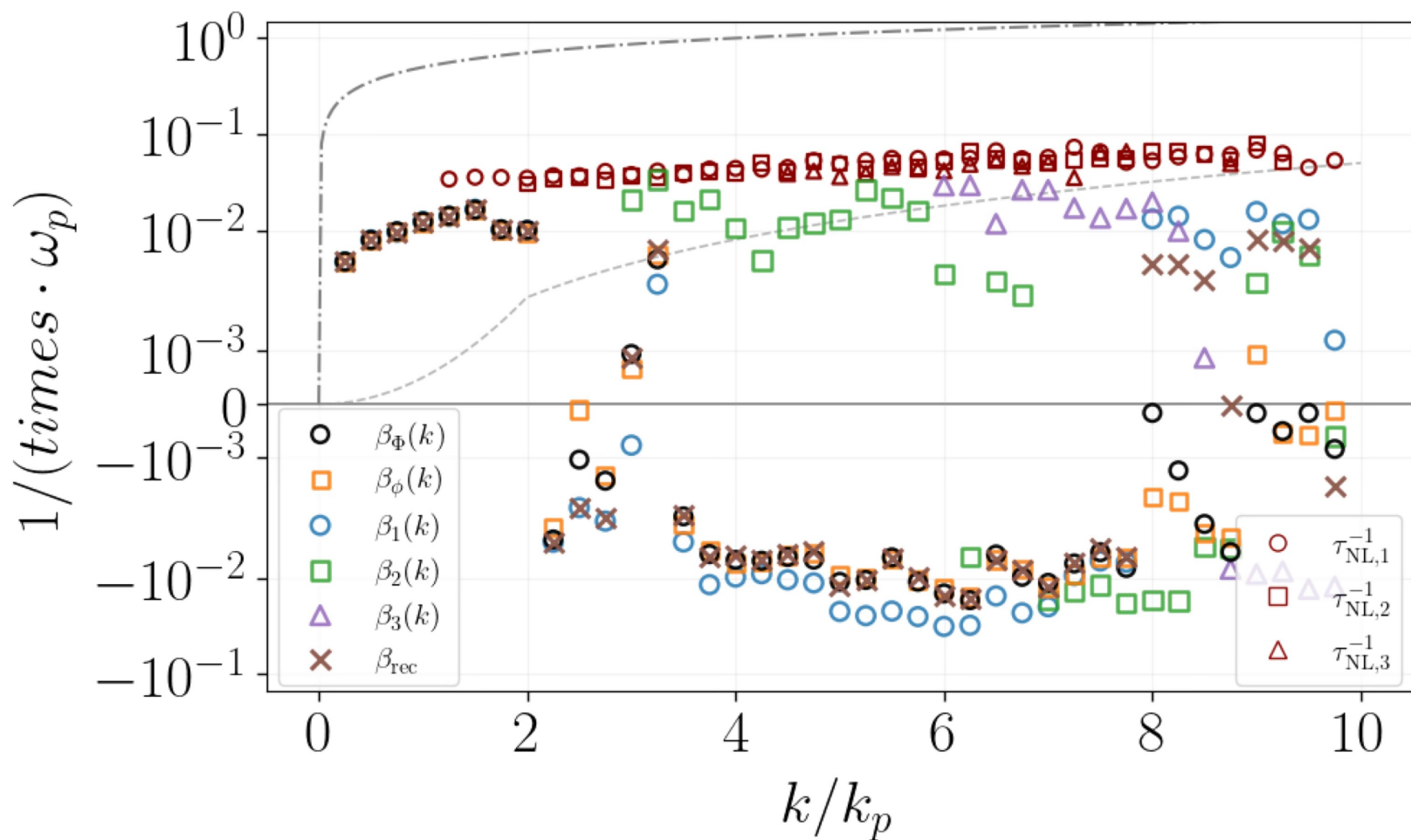
Energy Extraction: Growth rate per branch accounting for Doppler shift $\beta_N(k, N)$



Growth rate comparison $\beta_N(k, N)$

Wave growth rate $\beta_\phi(k), \beta_\Phi(k), \beta_{rec}, \beta_N(k)$

Inverse time scales $\tau_{NL,N}, \tau_{diss}, \tau_L$



Parserval theorem

$$\int \sum_N \Phi_N(k, t) dk = \int \Phi(k, t) dk = \int \phi(k, t) dk = \text{var}(\eta)$$

